LINEAR SORPTION MODELLED BY A MARKOV PROCESS

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Abstract

Consider a network of n compartments assuming that transition processes may occur between them: e.g. heat or mass transfer, or any kind of exchange of the first order. In fact, we have a Markov process with n states.

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1. Introduction

Consider a network of n compartments assuming that transition processes may occur between them: e.g. heat or mass transfer, or any kind of exchange of the first order. In fact, we have a Markov process with n states: A_1, A_2, \ldots, A_n ; the state occupancies x_1, x_2, \ldots, x_n (e.g. the distribution of 1 kg mass) obey the differential equations:

$$\frac{\mathrm{d}}{\mathrm{d}t} \begin{bmatrix} x_1\\ \vdots\\ x_n \end{bmatrix} = \begin{bmatrix} k_{11} & \dots & k_{1n}\\ k_{21} & \dots & k_{2n}\\ \vdots\\ k_{n1} & \dots & k_{nn} \end{bmatrix} \begin{bmatrix} x_1\\ \vdots\\ x_n \end{bmatrix}; \qquad k_{ij} \ge 0, \qquad i \ne j \qquad (1.1)$$

with

$$k_{ii} = -\sum_{j \neq i} k_{ji} \tag{1.2}$$

and the initial conditions [1]:

for t = 0: $x_1 = 1$, $x_2 = \ldots = x_n = 0$. (1.3)

In order to handle the problem we introduce the Laplace transforms

$$\mathcal{L} \{ x_i(t) \} = \int_0^\infty e^{-st} x_i(t) dt \equiv x_i(s); \quad i = 1, 2, \dots, n.$$
 (1.4)

Let us consider the partitioning (without any practical meaning for the time being)

 A_1 : 'mobile' phase and $\{A_2, \ldots, A_n\}$: 'fixed' phase

and let us calculate the ratio

$$\frac{x(s)}{x_1(s)} \equiv J(s), \qquad (1.5)$$

where

$$x(s) = \sum_{i=1}^{n} x_i(s).$$
 (1.6)

Using the rules

$$\mathcal{L}\left\{\frac{\mathrm{d}x}{\mathrm{d}t}\right\} = sx(s) - x(t)\Big|_{t=0},\qquad(1.7)$$

and

$$\mathcal{L}\{1\} = \frac{1}{s} \tag{1.8}$$

and applying Cramer's rule we immediately have from the transformed solution of (1.1) - (1.2) - (1.3):

$$J(s) = \frac{|sI - K|}{s |sI^1 - K^1|},$$
(1.9)

where $K = [k_{ij}]$ and K^1 is obtained by deleting the first row and first column in K (and similarly I^1 , too).

What we are interested in is whether or not J remains unchanged (invariant) by changing our model to a moving-phase one: the system being situated in a CSTR, or in a one-dimensional tube where convectionaldiffusional transport takes place.

2. The Lumped-parameter System (CSTR)

Let us define the flux (with respect to phase A_1 , which is now moving in fact):

$$y(t) = \omega x_1(t) \left[\text{kg/s} \right], \qquad (2.1)$$

with

$$\omega: \quad \text{flow rate}\left[1/s\right]. \tag{2.2}$$

There holds the following

THEOREM 1: J, see (1.5), is independent of ω , i.e. J is an invariant [2] and can be calculated from (1.9).

As an application of this statement we determine the RTD (residence time distribution) of the tracer particles in the CSTR. The balance equation becomes

$$\frac{\mathrm{d}x}{\mathrm{d}t} + y = 0, \qquad \text{i.e.} \qquad \frac{\mathrm{d}x}{\mathrm{d}t} + \omega x_1 = 0. \tag{2.3}$$

Using the rule (1.7) we have

$$sx(s) + \omega x_1(s) = 1.$$
 (2.4)

In view of (1.5) and Theorem 1, we immediately obtain

$$(J(s)s+\omega) x_1(s) = 1,$$

i.e.

$$y(s) = \omega x_1(s) = \frac{\omega}{J(s)s + \omega} = \frac{\omega}{s + \omega}\Big|_{s = J(s)s}.$$
 (2.5)

We realize that if n = 1: J = 1 and denoting the flux in that case by y^1 , we obtain the simple rule

$$y(s) = y^{1}(s)\Big|_{s=J(s)s}$$
 (2.6)

3. The Higher Moments of RTD

Obviously y(s) in (2.6) is just the transformed density function of the RTD. Let us denote the first three so-called cumulants or semi-invariants (mean, variance and 3rd central moment) by

$$\kappa_1, \kappa_2 \quad \text{and} \quad \kappa_3 \tag{3.1}$$

and in the case of only one phase, by

$$K_1, K_2 \quad \text{and} \quad K_3, \tag{3.2}$$

respectively. Rewriting J(s) as

$$J(s) = Cf(s) + 1$$
 with $f(0) = 1$, (3.3)

where C is the equilibrium constant between the 'fixed' and 'mobile' phase, we finally obtain the general formulae [3, 2]:

$$\kappa_1 = (C+1)K_1, \tag{3.4}$$

$$\kappa_2 = (C+1)^2 K_2 - 2Cf'(0)K_1, \qquad (3.5)$$

$$\kappa_3 = (C+1)^3 K_3 - 6C(C+1)f'(0)K_2 + 3Cf''(0)K_1.$$
(3.6)

4. The Distributed Parameter System (Tube)

Let us define the flux as (the analogue of (2.1)):

$$y(z,t) = Lx_1(z,t),$$
 (4.1)

where z is the length coordinate [m] and L is a linear operator, e.g.

$$L = u - D\frac{\partial}{\partial z}, \qquad (4.2)$$

with

$$u:$$
 flow velocity $[m/s]$ (4.3)

and

$$D:$$
 diffusivity $[m^2/s]$. (4.4)

THEOREM 2:

$$\frac{x(z,s)}{x_1(z,s)} = J(s), \qquad (4.5)$$

with J(s) as defined in (1.9). Once again, the left-hand side is independent of z and u as well as D (invariant).

As an application of this theorem [4] we determine the respective RTD. The balance equation becomes

$$\frac{\partial x}{\partial t} + \frac{\partial y}{\partial z} = 0 \tag{4.6}$$

and in the Laplace domain we have because of (4.5)

$$\left(J(s)s + \frac{\partial}{\partial z}L\right)x_1(z,s) = x_1(z,t)\Big|_{t=0}$$
: prescribed (4.7)

whose solution (depending on the boundary conditions) is to be substituted in the transformed equation (4.1):

$$y(z,s) = Lx_1(z,s).$$
 (4.8)

However, we realize that if n = 1: J = 1, and denoting the flux in that case by y^1 , we have

$$y(z,s) = y^{1}(z,s)\Big|_{s=J(s)s}$$
 (4.9)

which is the counterpart of (2.6). Once again, y(z, s) is the transformed density function of RTD. For the first three cumulants the general relations (3.4) - (3.5) - (3.6) are valid.

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References

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Appendix

1. The higher moments of RTD for a CSTR without sorption are given by the simple formula:

$$\mu_m = \frac{m!}{\omega^m}; \qquad m = 1, 2, \dots$$
(A.1)

where

$$\mu_m = \int_{0}^{\infty} t^m y^1(t) dt; \qquad m = 1, \, , 2, \, \dots \tag{A.2}$$

Consequently the first 3 semi-invariants become

$$K_{1} = \mu_{1} = \frac{1}{\omega}, \qquad K_{2} = \int_{0}^{\infty} (t - \mu_{1})^{2} y^{1} dt = \frac{1}{\omega^{2}},$$

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$$K_3 = \int_0^{1} (t - \mu_1)^3 y^1 dt = \frac{2}{\omega^3}.$$
 (A.3)

Inserting these values in (3.4) - (3.5) - (3.6) we obtain at once the respective semi-invariants for a CSTR with an imbedded Markov process.

2. The first 3 semi-invariants of RTD in a continuous-flow one-dimensional tube reactor with diffusion are given by the simple formulae [3]:

$$K_1 = \frac{1}{u}z$$
, $K_2 = \frac{2D}{u^3}z$, $K_3 = \frac{12D^2}{u^5}z$. (A.4)

Once again, substituting these values in (3.4) - (3.5) - (3.6) we have the respective semi-invariants for the generalized linear chromatography [3].

3. Consider an arbitrary continuous-flow reactor (either of the above examples) with two phases (mobile and fixed) between which a Markov process according to the system matrix

$$K = \begin{bmatrix} -k_{21} & k_{12} \\ k_{21} & -k_{12} \end{bmatrix}$$
(A.5)

takes place. Then, in view of (1.9) and (3.3) we readily obtain

$$J(s) = \frac{s + k_{12} + k_{21}}{s + k_{12}}, \qquad f(s) = \frac{k_{12}}{s + k_{12}}$$
(A.6)

and

$$C = \frac{k_{21}}{k_{12}}, \qquad f'_0 = -\frac{1}{k_{12}}, \qquad f''_0 = \frac{2}{k_{12}^2}.$$
 (A.7)

4. The interested reader might put down the respective formulae (3.4) - (3.5) - (3.6) in the case of the imbedded Markov process for both types of continuous-flow reactors.

Notation

- f see Eq. (3.3)
- s Laplace transform variable, s^{-1}
- t time, s
- u flow velocity, m/s
- x the total state occupancy (in the mobile plus fixed phase)
- x_1 state occupancy in the mobile phase, kg (CSTR) or kg/m (tube)
- y flux, kg/s
- z length coordinate, m
- C equilibrium constant, see Eq. (3.3)
- D diffusivity, m²/s
- J see Eq. (1.5)
- K system matrix, see Eq. (1.1)
- K_i semi invariants in a system consisting only of one phase (i = 1, 2, 3)
- κ_i semi invariants (i = 1, 2, 3)

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\omega flow rate, s<sup>-1</sup>
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