A Distributed Plasticity Approach for Steel Frames Analysis Including Strain Hardening Effects

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Abstract

This paper focuses on the creation and numerical application of physically nonlinear plane steel frames analysis problems. The frames are analysed using finite elements with axial and bending deformations taken into account. Two nonlinear physical models are used and compared – linear hardening and ideal elastic-plastic. In the first model, distributions of plastic deformations along the elements and across the sections are taken into account. The proposed method allows for an exact determination of the stress-strain state of a rectangular section subjected to an arbitrary combination of bending moment and axial force. Development of plastic deformations in time and distribution along the length of elements are determined by dividing the structure (and loading) into the parts (increments) and determining the reduced modulus of elasticity for every part. The plastic hinge concept is used for the analysis based on the ideal elastic-plastic model. The created calculation algorithms have been fully implemented in a computer program. The numerical results of the two problems are presented in detail. Besides the stress-strain analysis, the described examples demonstrate how the accuracy of the results depends on the number of finite elements, on the number of load increments and on the physical material model. COMSOL finite element analysis software was used to compare the presented 1D FEM methodology to the 3D FEM mesh model analysis.

reduced modulus of elasticity, distributed plasticity, incremental method, linear hardening, plastic deformations

1 Introduction

The material hardening phenomenon in steel structures has been widely discussed in theoretical and practical experimental based research [1, 2]. It is commonly accepted that a piecewise linear hardening model properly represents an actual steel stress-strain relationship [3]. Plastic deformations in a structure, independent of a nonlinear material model, can be evaluated using several different approaches: using the concentrated plastic hinge theory [4, 5], defining semi-rigid connections [6], using the distributed plasticity approach [7] or by linearizing the nonlinear stress diagram of a cross section [8]. In most cases, especially when plasticity is 'concentrated' at the nodes, quite strict assumptions are made, which makes calculations relatively simple, nevertheless, the reliability of results may be insufficient.

In this paper the authors continue their previous research [17] on the non-linear stress-strain state in a cross section and its numerical application for steel frame analysis. The main goal of the current research is to evaluate the distribution of plastic deformations along the length and in depth (across the sections) of the elements by dividing the structure into multiple finite elements and assigning them different moduli of elasticity in case of material linear hardening effect. The suggested methodology is new compared to previous researches [4-8] because it is based only on the fundamental equilibrium and compatibility equations, i.e. an equilibrium between internal and external forces is satisfied in any point of structure at any given time and the plane section assumption is valid. Because no other simplifications are made in the model and an incremental method [9] is used for the analysis of the problem, the computations are relatively complex (the computational cost of calculations is high) even in the case of a simple plane frame with uniaxial stress state in the cross sections. Nevertheless, authors suggest that this method allows for theoretically exact solution of linearly hardening bar structures. The authors chose to implement these calcu-lations in the Matlab environment [10], which will later allow them to incorporate it in the structural optimization problems [11–14] and [24–26] (which is the global goal of this research). This paper presents two numerical examples showing how the solutions are influenced by the number of finite elements and load increments used in the analysis. At the end, optimal values for these variables are suggested. In addition, the results of two material models (i.e. linear hardening and ideal elastic plastic) are compared.

2 Main assumptions and physical material models

The numerical algorithms for the analysis of steel frames in this paper are based on the following assumptions:

- plane sections remain plane after bending;
- plastic deformations are distributed along the length of the structural elements (in the case of the linear hardening model);
- a plastic hinge can be formed in any cross section (in the case of ideal elastic plasticity);
- the influence of bending moments and axial forces are taken into account in both – elastic and plastic material work stages;
- the influence of shear forces and tangential stress in the cross sections are ignored;
- the stability [15] of the structures and dynamic effects of loads are not considered
- deformations are small, i.e. equilibrium and compatibility equations are written for an undeformed structure;

Both the physical material models that are used in this paper can be shown in one stress-strain diagram (Fig. 1), which has three deformation paths – elastic (with elastic modulus E), elastic-plastic (with E_h), and ideally plastic (if $E_h = 0$).

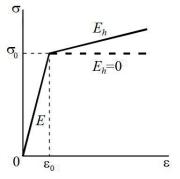


Fig. 1 Stress-strain diagrams of the linear hardening and ideal elastic plastic models

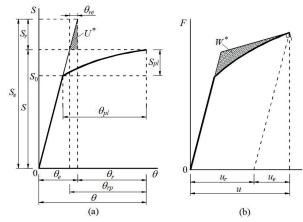


Fig. 2 (a) Relations between the internal forces and deformations and (b) between the loads and displacements

Stresses and strains of this diagram are related by the following equations:

$$E = \frac{\sigma_0}{\varepsilon_0}$$
 and $E_h = \frac{\sigma - \sigma_0}{\varepsilon - \varepsilon_0}$, (1)

 E_h – the modulus of elasticity in the second stage (linear hardening or ideal plasticity); σ_0 is the yield stress and ε_0 is the yield strain.

The elastic-plastic stage (when $\sigma > \sigma_0$) may be expressed in terms of the hardening ratio α :

$$\sigma = \sigma_0 + \alpha E \left(\varepsilon - \varepsilon_0 \right), \tag{2}$$

where $\alpha = E_{h}/E$ is the ratio between hardening and elastic moduli.

3 Main equations and general problem formulation

The stress state of a linearly hardening material is described by size *n* vectors of the total internal forces S, plastic internal forces S_{pl} , limit forces S_0 , residual forces S_r and elastic forces S_e (Fig. 2a). In the elastic-plastic system, vectors S, S_e and S_r are related as follows: $S = S_e + S_r$.

The deformed state of a structure is defined by size *m* vectors of the total displacements \boldsymbol{u}_r , residual displacements \boldsymbol{u}_r and elastic displacements \boldsymbol{u}_e (Fig. 2b), and size *n* vectors of the total deformations $\boldsymbol{\theta}_r$ residual deformations $\boldsymbol{\theta}_r$ and elastic deformations $\boldsymbol{\theta}_e$. These quantities are related as follows: $\boldsymbol{u} = \boldsymbol{u}_r + \boldsymbol{u}_e$; $\boldsymbol{\theta} = \boldsymbol{\theta}_r + \boldsymbol{\theta}_e$.

In addition, the residual deformations θ_r are divided into two parts: the residual elastic θ_{re} and residual plastic θ_{rp} deformations, thus $\theta_r = \theta_{re} + \theta_{rp}$. The residual internal forces are self-equilibrated, i.e. they satisfy the equilibrium equations:

$$AS_r = 0 , (3)$$

where A – is the equilibrium matrix ($m \times n$). The elastic internal forces are directly related to the external forces (loads) $F: AS_e = F$. The residual displacements and residual deformations are compatible, i.e. they satisfy the compatibility equations: $A^T u_r = \theta_r$.

In the elastic-plastic stress-strain states of the system, the internal forces and deformations are related through the flexibility matrix D and hardening matrix H:

$$\boldsymbol{D}\boldsymbol{S}_{r} = \boldsymbol{\theta}_{re}, \, \boldsymbol{H}\boldsymbol{\theta}_{rp} = \boldsymbol{S}_{pl}. \tag{4}$$

The elastic solution is determined from:

$$\boldsymbol{S}_{\boldsymbol{\rho}} = \boldsymbol{D}^{-1}\boldsymbol{A}^{T}(\boldsymbol{A}\boldsymbol{D}^{-1}\boldsymbol{A}^{T})^{-1}\boldsymbol{F} .$$
⁽⁵⁾

The yield conditions for a structure may be written as follows:

$$\boldsymbol{S}_{0} + \boldsymbol{H}(\boldsymbol{\theta}_{rp}^{+} + \boldsymbol{\theta}_{rp}^{-}) - \boldsymbol{\Phi}(\boldsymbol{S}_{e}^{-} + \boldsymbol{S}_{r}^{-}) \ge \boldsymbol{0} , \qquad (6)$$

where $\boldsymbol{\Phi}$ is the yield matrix; $\boldsymbol{\theta}_{rp}^+$ is the vector of the only positive values of vector $\boldsymbol{\theta}_{rp}$ (non-positive values of vector $\boldsymbol{\theta}_{rp}$ are equated to zero); $\boldsymbol{\theta}_{rp}^-$ is the vector of the only negative values of vector $\boldsymbol{\theta}_{rp}$ with the opposite signs (the negative signs are changed to positive and initially the non-negative values of vector $\boldsymbol{\theta}_{rp}$ are equated to zero). The last-mentioned vectors are related: $\boldsymbol{\theta}_{rp} = \boldsymbol{\theta}_{rp}^+ - \boldsymbol{\theta}_{rp}^-$. Together, the Eqs. (3) and (6) define the domain of statically admissible solutions of the residual internal forces.

Static formulation of the problem, in the case when the relation between plastic internal forces and deformations is linear, was introduced by Čyras [16]. In this case, the classical extreme mechanics principle is used: *of all statically admissible vectors of residual internal forces, the actual one corresponds to the minimum sum of complementary and potential plastic deformation energies.*

The complementary deformation energy U^* (Fig. 2a) is expressed as follows:

$$U^* = \int_{\theta_e - \theta_{re}}^{\theta_e} \mathbf{S}_r d\theta = \frac{1}{2} \mathbf{S}_r^T \mathbf{D} \mathbf{S}_r.$$
(7)

In this paper the relation between plastic internal forces and deformations is nonlinear, therefore, the potential plastic deformation energy is expressed by the integral:

$$U_{pl} = \int_{\theta-\theta_{pl}}^{\theta} \boldsymbol{S}_{pl} d\theta = \int_{\theta-\theta_{pl}}^{\theta} \left(\theta_{rp}^{+T} - \theta_{rp}^{-T}\right) \boldsymbol{H} d\theta .$$
(8)

Then the mathematical model of the problem stated on the basis of the above-mentioned principle reads:

$$\left(\frac{1}{2}\boldsymbol{S}_{r}^{T}\boldsymbol{D}\boldsymbol{S} + \int_{\boldsymbol{\theta}-\boldsymbol{\theta}_{pl}}^{\boldsymbol{\theta}} \left(\boldsymbol{\theta}_{p}^{+T} - \boldsymbol{\theta}_{rp}^{-T}\right) \boldsymbol{H} d\boldsymbol{\theta}\right) \to \min, \qquad (9)$$

$$\boldsymbol{S}_{0} + \boldsymbol{H} \left(\boldsymbol{\theta}_{rp}^{+} + \boldsymbol{\theta}_{rp} \right) - \boldsymbol{\Phi} \left(\boldsymbol{S}_{e}^{+} + \boldsymbol{S}_{r} \right) \ge \boldsymbol{0}, \quad \boldsymbol{A}\boldsymbol{S}_{r}^{-} = \boldsymbol{0}, \quad (10)$$

$$\boldsymbol{\theta}_{rp}^{+} \ge \mathbf{0}, \ \boldsymbol{\theta}_{rp}^{-} \ge \mathbf{0} \ , \tag{11}$$

The dual (kinematic) problem formulation reads:

$$\left(-\frac{1}{2} \boldsymbol{S}_{r}^{T} \boldsymbol{D} \boldsymbol{S}_{r} + \int_{\boldsymbol{\theta}-\boldsymbol{\theta}_{pl}}^{\boldsymbol{\theta}} \left(\boldsymbol{\theta}_{rp}^{+T} - \boldsymbol{\theta}_{rp}^{-T} \right) \boldsymbol{H} d\boldsymbol{\theta} - \left(\boldsymbol{\theta}_{rp}^{+T} - \boldsymbol{\theta}_{rp}^{-T} \right) \boldsymbol{H} \left(\boldsymbol{\theta}_{rp}^{+} - \boldsymbol{\theta}_{rp}^{-} \right) + \left(\boldsymbol{\theta}_{rp}^{+T} - \boldsymbol{\theta}_{rp}^{-T} \right) \left(\boldsymbol{\Phi} \boldsymbol{S}_{e} - \boldsymbol{S}_{0} \right) \right) \rightarrow \max,$$

$$(12)$$

$$\boldsymbol{D}\boldsymbol{S}_{r} + \boldsymbol{\Phi}^{T} \left(\boldsymbol{\theta}_{rp}^{+T} - \boldsymbol{\theta}_{rp}^{-T}\right) - \boldsymbol{A}^{T} \boldsymbol{u}_{r} = \boldsymbol{0},$$
(13)

$$\boldsymbol{\theta}_{rp}^{+} \geq \mathbf{0}; \, \boldsymbol{\theta}_{rp}^{-} \geq \mathbf{0} \,. \tag{14}$$

The constraints (13)–(14) define the kinematically admissible distributions of the residual deformations and displacements. Equations (13) actually denote the compatibility between the residual deformations θ_r and its components θ_{re} and θ_{rv} .

It can be demonstrated that the objective function of the kinematic formulation (12) expresses the complementary work W^* of the external loads (Fig. 2b). Thus, the problem (12)–(14) corresponds to the extreme energy principle: *of all kinematically admissible vectors of residual displacements, the actual one corresponds to the maximum complementary work of external loads*.

4 Stress-strain state in a cross section of an element under bending and tension or compression4.1 Reduced modulus of elasticity

When physical nonlinearity is considered, the size and distribution of plastic deformations in a cross section is characterized by the multiple different parameters (the material properties E_h and E; strains; elastic and plastic normal stresses; heights of plastic zones). It is convenient to define only one variable that would define all these parameters in an otherwise very complex analysis problem. In the current methodology, such a variable is E_r – the reduced modulus of elasticity, which allows for the evaluating of the influence of E_h and E altogether (Fig. 3a). In replacing a cross section having E_h and E (Fig. 3b) with a section having the generalized E_r (Fig. 3c), it is imperative to ensure that the distribution of the strains ε remains the same.

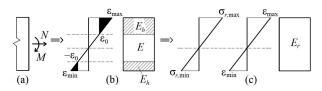


Fig. 3 (a) Internal forces in a cross section; (b) longitudinal strains, when two different moduli, E_h and E are considered; (c) normal stresses and longitudinal strains when one elastic modulus E_r is considered

Therefore, the following equality must be satisfied:

$$E_r = \frac{\sigma_{r,\max} - \sigma_{r,\min}}{\varepsilon_{\max} - \varepsilon_{\min}},$$
(15)

where $\sigma_{r,\max}$, $\sigma_{r,\min}$ are the values of normal stresses at the top and bottom of the cross section with the uniform elastic modulus (Fig. 3c); ε_{\max} , ε_{\min} are the longitudinal strains at the top and the bottom of the cross section (they have the same values in Fig. 3b and c).

4.2 Stress-strain state

The internal forces M and N should be in equilibrium with the normal stresses σ at every cross section:

$$N = \int_{A} \sigma(y) dA, \quad M = \int_{A} y \sigma(y) dA, \quad (16)$$

where A is the cross section area; y – the distance to the neutral axis of the cross section.

In addition, according to the classical mechanic's assumption of plane sections, every cross section must satisfy the equation:

$$\kappa = -\frac{\varepsilon_{\max} - \varepsilon_{\min}}{h}, \qquad (17)$$

where κ is the curvature and *h* is the height of the cross section. Thus, using the Eqs. (2), (16) and (17) we can write the system of equations fully describing the relations between stresses, strains and internal forces. In this paper we will only show the stress-strain state definition for the

case when both the bending moment and axial force are positive (the positive directions are shown in Fig. 3a). For example, for a rectangular cross section if normal stress is considered positive at the top $\sigma_{\text{max}} < \sigma_0$ and negative at the bottom $-\sigma_0 < \sigma_{\text{min}} < 0$ (III case in Fig. 4a), and $N < N_{\text{lim3}}$ (where N_{lim3} is the third axial force limit, which together with the bending moment causes zero normal stress at the bottom of the cross section $\sigma_{\text{min}} = 0$) there are seven equations to be written:

$$0.5bh_{el1}\sigma_{\min} + 0.5b\sigma_{0}h_{el2} + b\sigma_{0}h_{pl} + 0.5b(\sigma_{\max} - \sigma_{0})h_{pl} = N,$$
(18)

$$-\frac{1}{3}b\sigma_{\min}h_{el1}^{2} + \frac{1}{3}b\sigma_{0}h_{el2}^{2} + b\sigma_{0}h_{pl}\left(0.5h_{pl} + h_{el2}\right) +$$

$$0.5b(\sigma_{\max} - \sigma_{0})h_{pl}\left(\frac{2}{3}h_{pl} + h_{el2}\right) = M - Ny_{0},$$
(19)

$$\sigma_0 - \alpha \sigma_0 + E_h \kappa (y_0 - 0.5h) = \sigma_{\max},$$

$$E\kappa (y_0 + 0.5h) = \sigma_{\min}, \quad -E\kappa h_{el2} = \sigma_0,$$
(20)

$$y_0 + 0.5h = h_{el1}, \ 0.5h - y_0 - h_{el2} = h_{pl}.$$
 (21)

In this system, Eq. (18) defines the equilibrium of forces perpendicular to the plane of the cross section; equation (19) defines the equilibrium of bending moments in respect to the neutral axis. Equations (20) relate stresses and curvature in different zones of the cross section (taking into account the 'plane sections' assumption (17) and Eq. (2)). Equations (21) relate to the heights of elastic deformations zones h_{el1} , h_{el2} , the height of plastic deformations zone h_{pl} and the ordinate of neutral axis y_0 . For a more detailed explanation of these equations refer to [17].

The seven equation system (18)–(21) can be mathematically simplified to:

$$\overline{a}\kappa^2 y_0^2 + \overline{b}\kappa^2 y_0 + \overline{c}\kappa^2 + \overline{d}\kappa y_0 + \overline{e}\kappa + \overline{f} = 0,$$
(22)

Table 1	Equations	for determining	four axial force	limits (Fig.	4b) of a	rectangular cross section
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Axial force limits	Nonlinear equation	Coefficients of the nonlinear equation
$N_{\rm lim1} = \frac{6(M_0 - M)}{h}$	-	-
$N_{\rm lim2} = \sigma_0 b \left(h - h_{el} + \alpha \left(\frac{h^2}{h_{el}} - 2h + h_{el} \right) \right)$	$\overline{a}h_{el}^3 + \overline{b}h_{el}^2 + \overline{c}h_{el} - \overline{d} = 0$	$\overline{a} = 2\sigma_0 b (1-\alpha); \overline{b} = 3\sigma_0 b h (\alpha - 1); \overline{c} = 6M;$ $\overline{d} = \alpha \sigma_0 b h^3$
$N_{\rm lim3} = \sigma_0 b \left(h - 0.5 h_{el} + \alpha \left(\frac{h^2}{2h_{el}} - h + 0.5 h_{el} \right) \right)$	-	$\overline{a} = \sigma_0 b (1-\alpha); \overline{b} = 1.5 \sigma_0 b h (\alpha - 1); \overline{c} = 6M;$ $\overline{d} = 0.5 \alpha \sigma_0 b h^3$
$N_{\rm lim4} = \sigma_0 bh + \frac{6M}{h}$	-	-

$$\overline{g}\kappa^{3}y_{0}^{3} + \overline{h}\kappa^{3}y_{0}^{2} + \overline{i}\kappa^{3}y_{0} + \overline{j}\kappa^{2}y_{0}^{2} + \overline{k}\kappa^{3} + \overline{l}\kappa^{2}y_{0} + \overline{m}\kappa^{2} + \overline{n} = 0,$$
(23)

where $\overline{a}, \overline{b}, \dots \overline{n}$ are the coefficients of the nonlinear equations (Table 5). All possible stress-strain states (see Table 5 in Appendix A and Fig. 4a) and axial force limits (Table 1 and Fig. 4b) are explained in detail in [17].

Table 2 allows the determination of which combination of stress-strain state variations is to be used for a particular cross section subjected to the bending moment and axial force. This table shows all possible combinations of the stress-strain states and axial force limits when the axial force increases and the bending moment remains constant. In advance of using Table 2, the axial force limits have to be calculated according to Table 1. Then, following the ascending order of the axial force limits' values and the value of the particular axial force under consideration, the necessary stress-strain state is determined from Table 2 and the corresponding system of equations – from Table 5 (Appendix A). For example, let's say that the bending moment *M* is smaller than M_0 and the equations of Table 1 yields two real values of N_{lim2} , three of N_{lim3} , one of N_{lim1} and one of N_{lim4} . Let's say that all these values and the value of the given axial force lie in the following order $N_{\text{lim1}} < N_{\text{lim3}}^a < N_{\text{lim2}}^b < N_{\text{lim2}}^b < N < N_{\text{lim3}}^c < N_{\text{lim4}}^c$. Table 2 indicates that this order corresponds to 'Combination VI' and the position of the axial force *N* indicates the use of the third stress-strain state equations for the analysis.

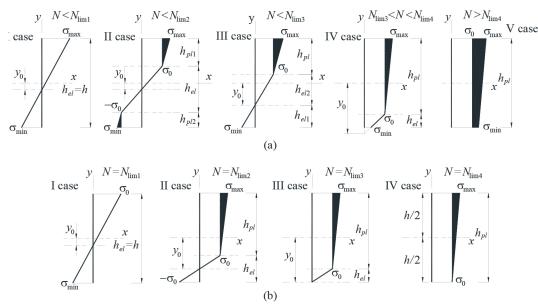


Fig. 4 Nine stress-strain states of a rectangular cross section when (a) $N \neq N_{\text{lim}}$ and when (b) $N = N_{\text{lim}}$

Table 2 Combinations of stress-strain states variations in a rectangula	ar cross section when the axial force is gradually increasing
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Combination	Stress-strain state	$N_{ m lim}$	Stress-strain state	$N_{ m lim}$	Stress-strain state	$N_{ m lim}$	Stress-strain state	$N_{ m lim}$	Stress-strain state	$N_{ m lim}$	Stress-strain state	$N_{ m lim}$	Stress-strain state	$N_{ m lim}$	Stress-strain state	Value of bending moment
Ι	1st	1	4th	4	5th	-	-	-	-	-	-	-	-	-	-	$M {<} M_0$
II	1st	1	3rd	3	4th	4	5th	-	-	-	-	-	-	-	-	$M \leq M_0$
III	1st	1	4th	3ª	3rd	3 ^b	4th	4	5th	-	-	-	-	-	-	$M \leq M_0$
IV	1st	1	4th	3ª	3rd	2ª	2nd	2 ^b	3rd	3 ^b	4th	4	5th	-	-	$M \leq M_0$
V	1st	1	3rd	3ª	4th	3 ^b	3rd	3°	4th	4	5th	-	-	-	-	$M \leq M_0$
VI	1st	1	3rd	3ª	4th	3 ^b	3rd	2ª	2nd	2 ^b	3rd	3°	4th	4	5th	$M \leq M_0$
VII	1st	1	3rd	2ª	2nd	2 ^b	3rd	3	4th	4	5th	-	-	-	-	$M \leq M_0$
VIII	-	2	3rd	3	4th	4	3rd	-	-	-	-	-	-	-	-	$M = M_0$
IX	-	2ª	3rd	2 ^b	2nd	2°	3rd	3	4th	4	5th	-	-	-	-	$M = M_0$
Х	2nd	2	3rd	3	4th	4	5th	-	-	-	-	-	-	-	-	$M > M_0$
XI	2nd	2ª	3rd	2 ^b	2nd	2°	3rd	3	4th	4	5th	-	-	-	-	$M > M_0$

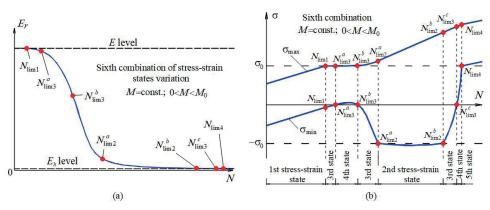


Fig. 5 (a) Typical variations of the reduced modulus of elasticity E_r and (b) normal stresses at the top layer σ_{max} and the bottom layer σ_{min} of a cross section while the axial force increases

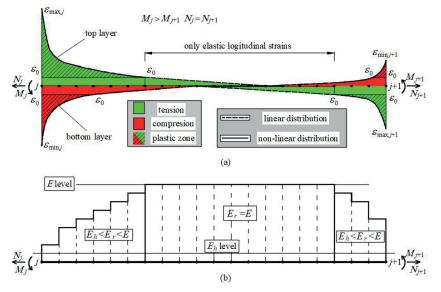


Fig. 6 A structural element divided into finite elements: (a) characteristic nonlinear distribution of longitudinal strains; (b) distribution of the reduced moduli of elasticity E_{μ}

Fig. 5 shows three variation curves of the unknowns when stress-strain states and axial force limits change according to 'Combination VI' of Table 2 (all curves are drawn from repetitive analysis when N increases and M = const.). It is worth mentioning that normal stress at the bottom of the cross section σ_{min} has a general tendency to decrease in a particular part of the loading curve although the axial force constantly increases (Fig. 5b).

5 Strain distribution along the length of an element

The procedure described in section 4 allows the determination of the stress-strain state at any point of an element k. A typical nonlinear distribution of longitudinal strains in the bottom and top layers of a cross section determined from internal forces at multiple points of an element is shown in Fig. 6. Distribution of the plastic strains along the length of a structural element can be determined with reasonable precision by dividing it into r_k finite elements. Naturally, the solution converges to an 'exact' one if $r_k \rightarrow \infty$. In the current methodology we evaluate the plastic zone of every cross section through the reduced modulus of elasticity E_r . A typical distribution of E_r among the chosen number of finite elements is shown in Fig. 6b.

6 Development of plastic strains

If the load increases in an elastic-plastic structure, then at some point plastic strains start to develop. If, from that instant, the load is increased by the increment ΔF , then the stiffness of the structure starts to change – it reduces. Therefore, in contrast to an elastic system, a one load iteration is not sufficient to solve an analysis problem when a physically nonlinear material is evaluated. The loading must be divided into increments [9] – usually equal parts – ΔF^f , f=1, 2, ..., t where t is the number of load increments. Naturally, the accuracy of the results increases if the load increments are smaller. During a one load increment the bending moment and axial force increase, but while the variation of the bending moment within the element is changing, the axial force usually remains constant over the element length.

A reduction of element stiffness can be expressed in terms of the reduced modulus of elasticity E_r , which (in plastic state) changes over the element length l_k as well as over the load increment ΔF^f , making the two-way nonlinear variation surface (Fig. 7a). In the proposed incremental analysis, it is important to determine the average value of this nonlinear surface as precisely as possible. For an approximate analysis the middle point value of E_r may be accurate enough (Fig. 7a):

$$E_{r,k}^f = E_{r\mathrm{mid},k}^{f'} \,. \tag{24}$$

However, the exact value of $E_{r,k}^{f}$ is slightly higher. According to the increase of E_r values over the element length (Fig. 7b) and over the load increment (Fig. 7c) $E_{r,k}^{f}$ may be expressed in terms of the nine perimeter values of E_r :

$$E_{r,k}^{f} = 0,25\psi^{2} \left(E_{r\max,k}^{(f-1)} + E_{r\min,k}^{(f-1)} + E_{r\max,k}^{f} + E_{r\min,k}^{f} \right) - 0,5\psi^{2} \left(E_{r\min,k}^{f} + E_{r\min,k}^{(f-1)} + E_{r\max,k}^{f'} + E_{r\min,k}^{f'} - 2E_{r\min,k}^{f'} \right) + (25) + 0,5\psi \left(E_{r\min,k}^{f} + E_{r\min,k}^{(f-1)} + E_{r\max,k}^{f'} + E_{r\min,k}^{f'} - 4E_{r\min,k}^{f'} \right) + E_{r\min,k}^{f'}.$$

The multiple numerical experiments performed by the authors of this paper for the determination of the value ψ showed that generally $\psi = 1/3$ and even for relatively complex functions $\psi = 1/3$. In the latter case, it was observed that if the interval of the perimeter values decreases, then ψ converges to an exact value of 1/3. Therefore, it can be generally assumed that if $l_k \rightarrow 0$ and $\Delta F \rightarrow 0$, then $\psi \rightarrow 1/3$, i.e. the accuracy of the solution increases if the numbers of the finite elements and load increments are increased. If $\psi = 0$, Eq. (25) becomes a simplified version for the approximate analysis: $E_{r,k}^{f} = E_{rmid,k}^{f'}$.

7 Mathematical models and solution algorithms 7.1 Linear hardening material model

Let the structure be additionally divided into $\overline{k} = 1, 2, ..., \overline{r}$ finite elements and the load F divided into a set of the equal increments ΔF^{f} ; f = 1, 2, ..., t. Then a classical mathematical model for elastic frame analysis consisting of equilibrium, compatibility and constitutive equations can be adapted for an analysis of a physically nonlinear frame:

$$A\Delta S^{f} = \Delta F^{f}, \qquad (26)$$

$$A^{T} \Delta \boldsymbol{u}^{f} - \boldsymbol{Z} \Delta \boldsymbol{\varepsilon}^{f} = \boldsymbol{0} , \qquad (27)$$

$$\boldsymbol{D}_{\varepsilon}^{f} \Delta \boldsymbol{S}^{f} - \Delta \varepsilon^{f} = \boldsymbol{0} , \quad f = 1, 2, \dots, t , \qquad (28)$$

where *A* is the equilibrium matrix; $Z = diag Z_{\bar{k}}$ and $D_{\varepsilon}^{f} = diag D_{\varepsilon,\bar{k}}^{f}$ are the block-diagonal matrices of an entire structure consisting of the individual blocks for every element $- Z_{\bar{k}}$ and $D_{\varepsilon,\bar{k}}^{f}$; ΔS^{f} , Δu^{f} and $\Delta \varepsilon^{f}$ are the vectors of the increments of internal forces, displacements and longitudinal strains. The system of Eqs. (26)–(28) can be expressed in terms of displacements:

$$\mathbf{K}^{f} \Delta \mathbf{u}^{f} = \Delta \mathbf{F}^{f} \tag{29}$$

where $\mathbf{K}^{f} = A(\mathbf{Z}\mathbf{D}_{\varepsilon}^{f})^{-1}A^{T}$ is the stiffness matrix of the structure during *f*-th load increment. The matrix $\mathbf{Z}_{\vec{k}}$ of an individual finite element defines the known relation between the increments of deformations $\Delta \theta_{\vec{k}}^{f}$ and increments of longitudinal strains $\Delta \varepsilon_{\vec{k}}^{f}$:

$$\Delta \theta_{\bar{k}}^{f} = \begin{vmatrix} \Delta \varphi_{\bar{k}1}^{f} \\ \Delta \varphi_{\bar{k}2}^{f} \\ \Delta \delta_{\bar{k}}^{f} \end{vmatrix} = \mathbf{Z}_{\bar{k}} \Delta \boldsymbol{\mu}_{\bar{k}}^{f} = \frac{l_{k} W_{k}}{6I_{k}} \boxed{\begin{array}{c} 2 & 1 \\ 1 & 2 \\ \hline & 6I_{k} / W_{k} \\ \hline & 6I_{k} / W_{k} \\ \hline \end{array}} \begin{vmatrix} \Delta \varepsilon_{\mathbf{M}_{\bar{k}2}}^{f} \\ \Delta \varepsilon_{\mathbf{N}_{\bar{k}}}^{f} \\ \Delta \varepsilon_{\mathbf{N}_{\bar{k}}}^{f} \\ \hline \end{array} \right|, (30)$$

where $\Delta \varepsilon_{M_{\bar{k}1}}^f$, $\Delta \varepsilon_{M_{\bar{k}2}}^f$ are the increments of maximum longitudinal strains due to the increments of bending moments in the corresponding element nodes; $\Delta \varepsilon_{N_{\bar{k}}}^f$ – the increment of

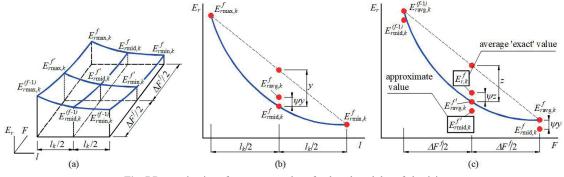


Fig. 7 Determination of an average value of reduced modulus of elasticity

maximum longitudinal strain due to the increment of axial force; $\Delta \varphi_{\bar{k}1}^f$, $\Delta \varphi_{\bar{k}2}^f$ – the increments of rotations in element nodes; $\Delta \delta_{\bar{k}}^f$ – the increment of element elongation. The matrix $\boldsymbol{D}_{\varepsilon,\bar{\kappa}}^f$ of every individual element depends

The matrix $D_{\varepsilon,\overline{k}}^{f}$ of every individual element depends on the stress state reached in the previous load increment. This matrix defines the known relation between the increments of longitudinal strains $\Delta \varepsilon_{\overline{k}}^{f}$ and increments of internal forces ΔS^{f} :

$$\Delta \boldsymbol{\mu}_{\overline{k}}^{f} = \boldsymbol{D}_{\varepsilon,\overline{k}}^{f} \Delta S_{\overline{k}}^{f} = \frac{1}{E_{r,\overline{k}}^{f} W_{\overline{k}}} \boxed{\begin{array}{c}1\\\\1\\\\\hline\end{array}} \underbrace{1\\\\W_{\overline{k}}/A_{\overline{k}}\end{array}} \cdot \frac{\Delta M_{\overline{k}1}^{f}}{\Delta M_{\overline{k}2}^{f}} (31)$$

where $E_{r,\bar{k}}^f$ is the reduced modulus of elasticity determined according to Eqs. (24) or (25). In every step of the iterative analysis process (Fig. 8) this modulus has to be determined for every plastically deformed finite element.

For example, let's say that for the *f*-th load increment – $E_{r,\overline{k}}^f = E_{r,\overline{k}}^{(f-1)}$. Every external iteration begins with the first internal iteration -p = 1, i.e. $E_{r,\overline{k}}^{f,p} = E_{r,\overline{k}}^{f,1}$ and problem (26)–(28) is solved. The internal forces determined in this internal iteration allows for the calculating of a new reduced modulus of elasticity $E_{r,\overline{k}}^{f,2}$ and the solving of problem (26)–(28) again. Internal iterations are continued while the relation $\left|E_{r,\overline{k}}^{f,p} - E_{r,\overline{k}}^{f,(p-1)}\right| / E_{r,\overline{k}}^{f,p}$ becomes smaller than the determined convergence tolerance λ . The last internal iteration number

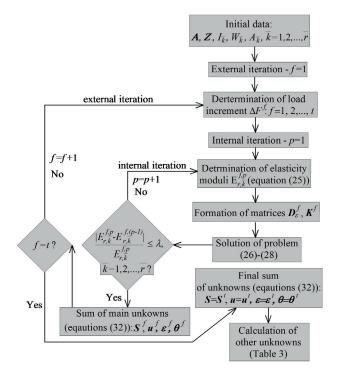


Fig. 8 Block scheme of analysis problem

is indicated as v, while the last value of E_r is carried to the first internal iteration of the next external iteration (f + 1). Fig. 9 shows the sequence of internal and external iterations used to determine E_r in the $F - u_i$ diagram.

When all load increments are evaluated, the total values of the unknowns are determined by combining the contributions of all *t* iterations:

$$\boldsymbol{u} = \sum_{f=1}^{t} \Delta \boldsymbol{u}^{f}; \quad \boldsymbol{S} = \sum_{f=1}^{t} \Delta \boldsymbol{S}^{f}; \quad \boldsymbol{\mu} = \sum_{f=1}^{t} \Delta \boldsymbol{\mu}^{f}; \quad \boldsymbol{\theta} = \sum_{f=1}^{t} \Delta \boldsymbol{\theta}^{f}. \quad (32)$$

If the total displacements, internal forces, strains and deformations are determined, all other parameters of the stress-strain state in the structure may be readily calculated. The equations for several more important parameters are given in Table 3.

7.2 Ideal elastic plastic material model

If all quantities associated with the plastic deformations θ_{rp}^+ and θ_{rp}^- are eliminated from model (9)–(11), it becomes a classical analysis problem formulation for a structure of elastic perfectly plastic material [21–22]:

$$\frac{1}{2}\boldsymbol{S}_{r}^{T}\boldsymbol{D}\boldsymbol{S}_{r} \to \min, \tag{33}$$

$$f(\boldsymbol{S}_e + \boldsymbol{S}_r) \le \boldsymbol{S}_0 , \qquad (34)$$

$$AS_r = \mathbf{0},\tag{35}$$

where the yield conditions (34) combines the nonlinear yield conditions of every section expressed in terms of the residual S_r and elastic S_e internal forces. For rectangular sections these conditions for every section are [21]:

$$\left|M_{e,j} + M_{r,j}\right| M_{0,j} + \frac{h_j^2}{16} \left|N_{e,j} + N_{r,j}\right|^2 \le M_{0,j}^2.$$
(36)

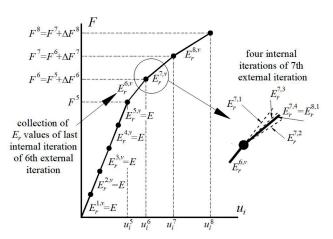


Fig. 9 Characteristic nonlinear force-displacement diagram and reduced modulus of elasticity E_r in external and internal iterations

Unknowns of a section j	Equations*
Maximum longitudinal strains (at the top and bottom of a section)	$\varepsilon_{\max,j} = \varepsilon_{M,j} + \varepsilon_{N,j} \; ; \; \; \varepsilon_{\min,j} = -\varepsilon_{M,j} + \varepsilon_{N,j}$
Maximum stress	$\sigma_{\max,j} = \sigma_0 + (\varepsilon_{\max,j} - \varepsilon_0)E_h; \ \sigma_{\min,j} = -\sigma_0 + (\varepsilon_{\min,j} - \varepsilon_0)E_h$
Heights of plastic and elastic zones	$\begin{split} h_{pl1,j} = h_j \frac{\varepsilon_{\max,j} - \varepsilon_0}{\varepsilon_{\max,j} - \varepsilon_{\min,j}}; h_{pl2,j} = h_j \frac{-\varepsilon_{\min,j} - \varepsilon_0}{\varepsilon_{\max,j} - \varepsilon_{\min,j}}; \\ h_{el,j} = h_j - h_{pl1,j} - h_{pl2,j} \end{split}$
Curvature	$\kappa_j = \frac{\varepsilon_{\min,j} - \varepsilon_{\max,j}}{h_j}$
Ordinate of the neutral axis	$y_{0,j} = h_j \left(0, 5 - \frac{\varepsilon_{\max,j}}{\varepsilon_{\max,j} - \varepsilon_{\min,j}} \right)$
Unknowns of a structure	Equations
Elastic and residual displacements	$\boldsymbol{u}_{e} = (\boldsymbol{A}\boldsymbol{D}^{-1}\boldsymbol{A}^{T})^{-1}\boldsymbol{F}; \ \boldsymbol{u}_{r} = \boldsymbol{u} - \boldsymbol{u}_{e}$
Elastic and residual internal forces	$\boldsymbol{S}_{e} = \boldsymbol{D}^{-1}\boldsymbol{A}^{T}\boldsymbol{u}_{e}; \ \boldsymbol{S}_{r} = \boldsymbol{S} - \boldsymbol{S}_{e}$

Table 3 Equations for determining the parameters of a structure and cross sections

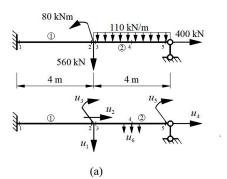
* these equations are valid only for the second limit state (see second scheme of Fig. 5a) if $\sigma_{\text{max}} > \sigma_0 > 0$ and $\sigma_{\text{min}} < -\sigma_0 < 0$. Equations for other states can be easily derived and are not shown here for brevity.

8 Numerical examples and discussions

All described algorithms were numerically implemented with MATLAB software.

8.1 Example 1

This example is dedicated to the numerical implementation of the analysis problem algorithm (Fig. 8) and investigation into how the number of finite elements and the number of force increments influence the accuracy of results. All force increments are chosen to be equal. In addition, the influence of the coefficient ψ to the accuracy of results is analyzed (if $\psi = 1/3$, calculations are 'exact' and if $\psi = 0$ – 'approximate'). The steel beam (Fig. 10a) of rectangular cross section is considered – $h \times b = 0.3$ m × 0.1 m, E = 205 GPa, $E_{h} = E/50$, $\sigma_{0} = 235$ MPa. Two types of finite elements are used: with and without distributed load [21,23]. Internal iterations of the analysis problem are stopped when the change of all E_r values is smaller than 0.1 % ($\lambda = 0.001$). Calculations have been performed for several values of the finite element number r_{μ} and load increments number t. It should be noted that the maximum value of stress reached in the beam ($\sigma_{\rm max} \cong 389~{\rm MPa}$, see Table 4) well exceeds the limit of the steel strength, in addition geometrical non-linearity is not evaluated and therefore this is a theoretical problem just to test the algorithm. The case of analysis, when $r_1 = r_2 = t = 256$ can be considered as the most accurate or almost absolutely accurate. However, the computational time to solve it is unacceptably long - it takes around 17 hours on an average PC. Fig. 12 and Table 4 shows that relatively good results (accuracy of \sim (1–2) %) can be achieved if r_k and t are equal to 16 or 32. Other significant results – the plastic internal forces, distributions of plastic zones – are shown in Figs. 10b and 11.



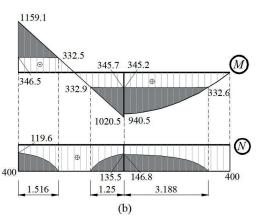


Fig. 10 (a) Scheme and discrete model of the beam; (b) total and plastic internal forces and plastic zones

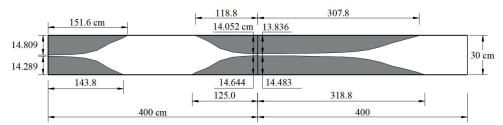


Fig. 11 Distribution of elastic and plastic zones in the beam

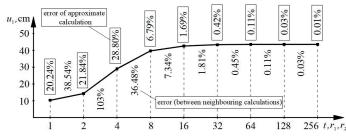


Fig. 12 Variation of displacement and error of displacement while simultaneously increasing the number of finite elements and load increments

	Table 4 Main results of the numerical example							
r_k/t	u_1 , cm	M_1 , kNm	$\sigma_{\rm max,l},$ MPa	u_1^*, cm	Calculation time			
1/1	10.3407	-1220.9721	250.6762	8.5998	<1 min (<1 min*)			
2/2	14.3255	-1267.4694	256.4700	17.4543	<1 min (<1 min*)			
4/4	29.0813	-1202.9674	295.2400	37.4555	<1 min (<1 min*)			
8/8	39.6896	-1166.8112	351.4875	42.3861	~3 min (<1 min*)			
16/16	42.6040	-1160.8554	372.7660	43.3246	~8 min (~1 min*)			
32/32	43.3741	-1159.5057	381.9948	43.5582	~23 min (~3 min*)			
64/64	43.5699	-1159.1653	386.3044	43.6161	~84 min (~9 min*)			
128/128	43.6189	-1159.0794	388.3873	43.6305	~270 min (~36 min*)			
256/256	43.6317	-1159.0622	389.4128	43.6346	~17 h (~150 min*)			

Explanations: ()* - 'approximate' solution (when $\psi = 0$); σ_{max} - maximum absolute value of normal stress in the cross section.

8.2 Example 2

In this example, the presented algorithm is applied for the analysis of a frame with linear hardening material model, and the results are compared to the solution of the elastic ideally plastic frame. A two story steel frame (Fig. 13a) with cross section parameters of $A_1 = 880 \text{ cm}^2 (0.44 \text{ m} \times$ 0.2 m), $A_2 = 480 \text{ cm}^2$ (0.4 m × 0.12 m), $A_3 = 300 \text{ cm}^2$ (0.3 m \times 0.1 m) and material properties E = 205 GPa and $\sigma_0 =$ 235 MPa is considered. The analysis problem is solved for two cases: 1) $E_{h} = 0$ (the material is elastic ideally plastic, – plasticity is concentrated at the nodes); 2) $E_{h} = E/50$ (the material is linearly hardening and close to real carbon steel properties, - plasticity is distributed). The analysis results are shown in Figs. 13b and 14. These Figs. indicate that the plastic zones and plastic hinges coincide. The node displacements of the ideally plastic solution are 44-51 % larger than the hardening material model solution $(E_{h} =$

E/50). For a qualitative comparison the same structure was modeled in COMSOL software using 3D FE and linear hardening material model (with $E_h = E/50$) (Fig. 14a). Determined maximum horizontal displacement (12.96cm) and plastic zone areas coincide with the results from the suggested methodology comparably well.

9 Conclusions

 An algorithm for linearly hardening plane frame analysis based only on classical stress-strain state assumptions in a cross section was described. The proposed algorithm allows for the evaluation of plastic deformations along the length of an element and across the height of a cross section. Self-correcting internal iterations in the algorithm allows for the determining of accurate solutions which can be used to qualify other approximate methods.

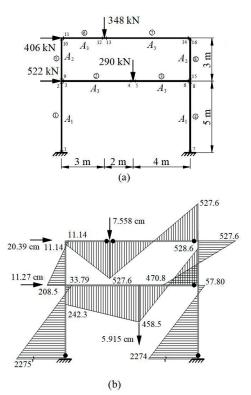


Fig. 13 (a) Frame model; (b) plastic hinges, bending moments and node displacements of the first analysis problem $(E_h = 0)$

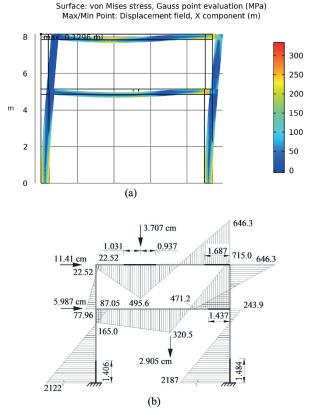


Fig. 14 (a) Deformed scheme, top beam displacement and von Mises stresses from COMSOL environment; (b) plastic zones, bending moments and node displacements of the second analysis problem

- 2. The proposed algorithm is suitable for implementation in a structural optimization problem, which is the current research topic of the paper's authors.
- 3. The numerical results (Example 1) indicate that it is sufficient to divide the structural elements into 16-32 finite elements and external loading into 16-32 increments to achieve reasonably accurate results. In addition, it shows that the 'accurate' solution method (when $\psi = 1/3$) very slightly improves the accuracy of the results compared to the 'approximate' solution (when $\psi = 0$): if $r_1 = r_2 = t = 32$, the values of displacement u_1 differs only by 0.42 % (Fig. 12). If the number of elements and/or load increments is further increased, the difference reduces even more. The computational time could be effectively reduced by applying the finite element grid of selective mesh.
- 4. The comparison of ideal elastic plastic $(E_h = 0)$ and the linear hardening material $(E_h = E/50)$ models shows that all the locations of the plastic deformations coincide, but the values of nodal displacements differ significantly: the first model gives 44–51 % bigger displacements. This difference is due to the plastic hinge concept: the hinges form instantly in an entire section, while the linear hardening model allows for a gradual shift from elastic to plastic state and thus the stiffness of a structure reduces at a slower rate. Verification problem solved with COMSOL program using 3D mesh model gives 13.58 % bigger top maximum horizontal displacement.
- 5. This publication describes a more academic rectangular steel section analysis, nevertheless the presented methodology together with the previously published analytical equations [17] can be applied for the elastic-plastic stress-strain analysis of any cross-section.

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Fifth stress-strain state	Fourth stress-s	train state	Third stress-strain state	Second stres	s-strain state	First stress- strain state	
	$\overline{a}\kappa^2 + \overline{b}\kappa + \overline{c} = 0$	$\overline{a}h_{el}^3 + \overline{b}h_{el}^2 - \overline{c}h_{el} + \overline{d} = 0$	$\overline{a}\kappa^2 y_0^2 + \overline{b}\kappa^2 y_0 + \overline{c}\kappa^2 + \overline{d}\kappa y_0 + +\overline{e}\kappa + \overline{f} = 0;$ $\overline{g}\kappa^3 y_0^3 + \overline{h}\kappa^3 y_0^2 + \overline{i}\kappa^3 y_0 + \overline{j}\kappa^2 y_0^2 + + \overline{k}\kappa^3 + \overline{l}\kappa^2 y_0 + \overline{m}\kappa^2 + \overline{n} = 0$	$\overline{a}h_{el}^3 + \overline{b}h_{el} + \overline{c} = 0$	$\overline{a}y_0^5 + \overline{b}y_0^4 + \overline{c}y_0^3 + \overline{d}y_0^2 + \overline{c}y_0 + \overline{f} = 0$		Nonlinear equations
	$\overline{\alpha} = EE_hbh_{el}\left(h^2 - hh_{el} + h_{el}^2/3\right) - 1/3\left(E_hEb^4 + E^2bh_{el}^3\right); \overline{c} = \sigma_0\left(N - \sigma_0bh\right);$ $\overline{b} = E\left(\sigma_0b\left(0.5h^2\left(1 + \alpha\right) - 0.5h_{el}^2\left(1 - \alpha\right) - hh_{el}\left(1 + \alpha\right)\right) - M + Nh_{el} - 0.5Nh\right)$	$ar{a} = (b/6)(E - E_h)(N - \sigma_0 bh); ar{b} = 0.25b(E - E_h)(6\sigma_0 W - 2M - Nh); \ ar{c} = ME_h bh; ar{d} = E_h W(3\sigma_0 W - 0.5Nh + 3M)$	$\begin{split} \overline{\alpha} &= 4Eb \left(E - E_{h} \right); \overline{b} = 4Ebh \left(E + E_{h} \right); \overline{c} = Ebh^{2} \left(E - E_{h} \right); \overline{d} = 8\sigma_{0} Eb \left(\alpha - 1 \right); \\ \overline{e} &= 4E \left(\sigma_{0} bh \left(1 - \alpha \right) - 2N \right); \overline{f} = 4\sigma_{0}^{2} b \left(1 - \alpha \right); \overline{g} = 8E^{2} b \left(E_{h} - E \right); \overline{h} = -12E^{2} bh \left(E + E_{h} \right); \\ \overline{i} &= 6E^{2} bh^{2} \left(E_{h} - E \right); \overline{f} = 12\sigma_{0} E^{2} b \left(1 - \alpha \right); \overline{k} = -E^{2} bh^{3} \left(E + E_{h} \right); \overline{i} = 12E^{2} \left(\sigma_{0} bh \left(\alpha - 1 \right) + 2N \right); \\ \overline{m} &= 3E^{2} \left(\sigma_{0} bh^{2} \left(1 - \alpha \right) - 8M \right); \overline{n} = 4\sigma_{0}^{3} b \left(\alpha - 1 \right) \end{split}$	$\overline{\alpha} = 12\sigma_0 b(\alpha - 1); \overline{b} = 0.25\sigma_0 b(1 - \alpha) \left(h^2 + 4y_0^2\right) - M + Ny_0;$ $\overline{c} = \alpha \sigma_0 b h \left(2y_0^2 + h^2/6\right)$	$\begin{split} \overline{\alpha} &= 48\sigma_0^3 b^3 \big(3\alpha^2 - 3\alpha - \alpha^3 + 1 \big); \overline{b} = 48N\sigma_0^2 b^2 \big(\alpha^2 - 2\alpha + 1 \big); \\ \overline{c} &= 48M\sigma_0^2 b^2 \big(\alpha^2 - 2\alpha + 1 \big) + 12\sigma_0 b N^2 \big(1 - \alpha \big) + 4\sigma_0^3 b^3 h^2 \big(3\alpha - 2\alpha^2 - 1 \big); \\ \overline{d} &= 48MN\sigma_0 b \big(1 - \alpha \big); \overline{c} = 3N^2 \big(4M + \sigma_0 b h^2 \big(1 - \alpha \big) \big); \overline{f} = N^3 h^2 \end{split}$,	Coefficients of nonlinear equations
$\begin{split} y_0 &= \left(\sigma_0 b h^3 \left(1 - 1.5 \alpha\right) - h \left(N h + 3M\right)\right) / 18M;\\ \kappa &= -\left(12M\right) / \left(E_h b h^3\right); \varepsilon_{\max} = \kappa \left(y_0 - 0.5h\right);\\ \varepsilon_{\min} &= \kappa \left(y_0 + 0.5h\right); \sigma_{\max} = \sigma_0 \left(1 - \alpha\right) + E_h \varepsilon_{\max};\\ \sigma_{\min} &= \sigma_0 \left(1 - \alpha\right) + E_h \varepsilon_{\min}; h_{el} = 0 \end{split}$	$y_{0} = \sigma_{0} / E\kappa + h_{el} - 0.5h;$ $h_{pl} = h - h_{el}; \varepsilon_{max} = \kappa (y_{0} - 0.5h);$ $\varepsilon_{min} = \kappa (y_{0} + 0.5h); \sigma_{max} = \sigma_{0} (1 - \alpha) + E_{h} \varepsilon_{max};$ $\sigma_{min} = \varepsilon_{min} E$ $y_{0} = (\sigma_{0} bh^{3} (1 - 1.5\alpha) - h(Nh + 3M)) / 18M;$ $\kappa = -(12M) / (E_{h} bh^{3}); \varepsilon_{max} = \kappa (y_{0} - 0.5h);$ $\varepsilon_{min} = \kappa (y_{0} + 0.5h); \sigma_{max} = \sigma_{0} (1 - \alpha) + E_{h} \varepsilon_{max};$		$egin{aligned} h_{el1} = 0.5h + y_0; h_{el2} = -\sigma_0 / (E\kappa); \ h_{pl} = h - h_{el1} - h_{el2}; arepsilon_{ ext{max}} = \kappa (y_0 - 0.5h); \ arepsilon_{ ext{min}} = \kappa (y_0 + 0.5h); \sigma_{ ext{max}} = \sigma_0 (1 - lpha) + E_h arepsilon_{ ext{max}}; \ \sigma_{ ext{min}} = arepsilon_{ ext{min}} E \end{aligned}$	$\begin{split} h_{pl1} &= 0.5h - y_0 - 0.5h_{el}; h_{pl2} = 0.5h + y_0 - 0.5h_{el}; \\ \kappa &= -(2\sigma_0)/(h_{el}E); \varepsilon_{\max} = \kappa (y_0 - 0.5h); \\ \varepsilon_{\min} &= \kappa (y_0 + 0.5h); \sigma_{\max} = \sigma_0 (1 - \alpha) + E_h \varepsilon_{\max}; \\ \sigma_{\min} &= \sigma_0 (\alpha - 1) + E_h \varepsilon_{\min} \end{split}$		$\begin{split} \sigma_{\max} &= (M/W) + (N/A); \sigma_{\min} = -(M/W) + (N/A); \\ \varepsilon_{\max} &= \sigma_{\max}/E; \varepsilon_{\min} = \sigma_{\min}/E; h_{el} = h; \\ \kappa &= (\sigma_{\min} - \sigma_{\max})/(hE); y_0 = \frac{\sigma_{\max} + \sigma_{\min}}{2E\kappa} \end{split}$	Equations for other unknowns

Appendix A