

BASIC FLOW DIFFERENTIAL EQUATIONS APPLIED IN BUILDING SERVICE ENGINEERING

László GARBAI and György SZÉKELY

Department of Building Engineering
Budapest University of Technology and Economics
H-1521 Budapest, Hungary
Phone: (36-1) 463 2405

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Abstract

In this article basic flow equations will be presented to describe the flow behaviour of various mediums, like gas, ideal gas, wet steam, ... etc. The equations constitute a system of simultaneous differential equations. The solution of this system is here not discussed. The Appendix gives a clearly arranged synopsis of the differential equations of flows under different conditions.

Keywords: district heating, gas supply, energy equation, flow equation.

1. Introduction

An essential knowledge of flow equations is required in process planning, like Building Service Engineering. There are several areas in practise where giving a detailed model description is needed. The main areas where flowing mediums are investigated are the followings

- gas supply
- water supply
- district heating
- central heating

Tests are often made of current taking place in

- diffusers and reducers
- valves and dampers
- nozzles

It is important for the designer and operating engineer to be able to determine as exactly as possible the changes in the thermo- and hydrodynamic parameters of the current medium. In our paper we will show how to write down the current simultaneous differential equation system for various current problems and how to solve those problems with the highest accuracy while using the simplest computer tools. In the studied apparatuses the currents can take place with mechanical losses,

either negligible mechanical and heat loss, or simply negligible heat loss. A current can take place in pipes or ducts with constant or variable cross sections.

The most common mediums are

- hot water
- wet steam
- dry saturated steam or superheated steam
- natural gas

These mediums' behaviour approaches the properties of the ideal gas or real gas.

In our paper we test the above descriptions of currents. We will present those differential equations with which we can determine, under stationary conditions, changes in the pressure, specific volume, and temperature of the current medium.

2. Basic Equations of Currents

The one-dimensional motion equation for stationary flows with friction:

$$w \frac{dw}{dz} = -\frac{1}{\varrho} \frac{dp}{dz} - \frac{\lambda}{2d} w^2.$$

The continuity equation:

$$A\varrho w = \text{const.}$$

The energy equation

$$\frac{d}{dz} \left(h + \frac{w^2}{2} \right) = \frac{\dot{q}}{\dot{m}}.$$

The equation of state for ideal and real gas:

$$pv = ZRT.$$

The condition equation for wet and dry saturated steam and the used approximations are:

$$\begin{array}{lllll} p & = & \delta_0 + \delta_1 T + \delta_2 T^2 + \dots & \rightsquigarrow & p \approx \delta_0 + \delta_1 T \\ h(v = \text{const.}, T) & = & b_0 + b_1 T + b_2 T^2 + \dots & \rightsquigarrow & h(v = \text{const.}, T) \approx b_0 + b_1 T \\ h(v, T = \text{const.}) & = & g_0 + g_1 v + g_2 v^2 + \dots & \rightsquigarrow & h(v, T = \text{const.}) \approx g_0 + g_1 T \\ v'' & = & \gamma_0 + \gamma_1 T + \gamma_2 T^2 + \dots & \rightsquigarrow & v'' \approx \gamma_0 + \gamma_1 T \end{array}$$

3. Special Cases of Flows

The special cases of flows are deduced from the differential equations of politropic case with the appropriate substitutions:

- for isothermal current $\frac{dT}{dz} = 0$
- for isochor current $\frac{dv}{dz} = 0$
- for isobar currents $\frac{dp}{dz} = 0$
- for adiabatic currents $\dot{q} = 0$
- for isentropic currents $\frac{dS}{dz} = 0$
- constant cross section $\frac{dA}{dz} = 0$
- without friction $\frac{\lambda}{2d} v \left(\frac{\dot{m}}{A}\right)^2 = 0$

4. The Solution of the System of Equation

For the solution of the system of equations we must specify the initial conditions. If $x = 0, p = p_0, T = T_0, h = h_0$ and $v = v_0$

$$\left(\frac{\partial h}{\partial T}\right)_{v_0} \approx b_1, \quad \left(\frac{\partial h}{\partial v}\right)_{T_0} \approx g_1, \quad \left(\frac{\partial p}{\partial T}\right)_{p_0} \approx \delta_1.$$

We need to know the functions

$$\begin{aligned} A &= A(z) \\ \dot{m} &= \text{const} \\ \dot{q} &= \dot{q}(z) \end{aligned}$$

The differential equations presented in the appendix can be solved by direct integration (after separation), by a method of successive approximation or with step-by-step finite difference method.

5. Summary

The tables summarize the main equations for special cases of flows. Engineers, scientists are often in need of having a table, which contains of the describing equations of flows and their solutions. In civil building service engineering it is useful to have the main equations collected.

SYMBOLS

T	temperature
w	velocity
v	specific volume
h	enthalpy
p	pressure
z	space coordinate
\dot{m}	mass flow
A	cross section
d	diameter
R	specific gas constant
ϱ	density
λ	coefficient of friction
\dot{q}	specific heat loss
U	perimeter
Z	compressibility factor
S	entropy
κ	adiabatic exponent

Table 1. Ideal Gas Flow Equations I.

Table 1. Ideal Gas Flow Equations II.

	Cross-Section	No Friction	No Heat Loss No Friction
Polytropic	Variable	$-\frac{\dot{m}^2}{A^3} v \frac{dA}{dz} = -\frac{R_t}{v} \frac{dT}{dz} - \left[\left(\frac{\dot{m}}{A} \right)^2 - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $\frac{\dot{q}}{m} + \dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz} = \frac{\kappa}{\kappa-1} R_t \frac{dT}{dz} + \frac{\dot{m}^2}{A^2} v \frac{dv}{dz}$	
	Constant	$0 = \frac{R_t}{v} \frac{dT}{dz} - \left[\left(\frac{\dot{m}}{A} \right)^2 - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $\frac{\dot{q}}{m} = \frac{\kappa}{\kappa-1} R_t \frac{dT}{dz} + \frac{\dot{m}^2}{A^2} v \frac{du}{dz}$	
Adiabatic	Variable		$-\frac{\dot{m}^2}{A^3} v \frac{dA}{dz} = -\frac{R_t}{v} \frac{dT}{dz} - \left[\left(\frac{\dot{m}}{A} \right)^2 - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $-\dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz} = \frac{\kappa}{\kappa-1} R_t \frac{dT}{dz} + \frac{\dot{m}^2}{A^2} v \frac{dv}{dz}$
	Constant		$0 = \frac{R_t}{v} \frac{dT}{dz} - \left[\left(\frac{\dot{m}}{A} \right)^2 - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $0 = \frac{\kappa}{\kappa-1} R_t \frac{dT}{dz} + \frac{\dot{m}^2}{A^2} v \frac{dv}{dz}$
Isentropic	Variable	$-\frac{\dot{m}^2}{A^3} v \frac{dA}{dz} = - \left[\frac{R_t}{v} \left(\frac{\partial T}{\partial v} \right)_S - \left(\frac{\dot{m}}{A} \right)^2 + \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $\frac{\dot{q}}{m} + \dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz} = \left[\frac{\kappa}{\kappa-1} R_t \left(\frac{\partial T}{\partial v} \right)_S + \frac{\dot{m}^2}{A^2} v \right] \frac{dv}{dz}$ $\left(\frac{\partial T}{\partial v} \right)_S = -(\kappa - 1) \frac{T}{v}$	$-\frac{\dot{m}^2}{A^3} v \frac{dA}{dz} = - \left[\frac{R_t}{v} \left(\frac{\partial T}{\partial v} \right)_S - \left(\frac{\dot{m}}{A} \right)^2 + \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $\dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz} = \left[\frac{\kappa}{\kappa-1} R_t \left(\frac{\partial T}{\partial v} \right)_S + \frac{\dot{m}^2}{A^2} v \right] \frac{dv}{dz}$
	Constant	$0 = \left[-\frac{R_t}{v} \left(\frac{\partial T}{\partial v} \right)_S - \left(\frac{\dot{m}}{A} \right)^2 + \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $\frac{\dot{q}}{m} = \left[\frac{\kappa}{\kappa-1} R_t \left(\frac{\partial T}{\partial v} \right)_S + \frac{\dot{m}^2}{A^2} v \right] \frac{dv}{dz}$ $\left(\frac{\partial T}{\partial v} \right)_S = -(\kappa - 1) \frac{T}{v}$	$0 = - \left[-\frac{R_t}{v} \left(\frac{\partial T}{\partial v} \right)_S - \left(\frac{\dot{m}}{A} \right)^2 + \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $0 = \left[\frac{\kappa}{\kappa-1} R_t \left(\frac{\partial T}{\partial v} \right)_S + \frac{\dot{m}^2}{A^2} v \right] \frac{dv}{dz}$ $\left(\frac{\partial T}{\partial v} \right)_S = -(\kappa - 1) \frac{T}{v}$
Isothermal	Variable	$-\frac{\dot{m}^2}{A^3} v \frac{dA}{dz} = - \left[\left(\frac{\dot{m}}{A} \right)^2 - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $\frac{\dot{q}}{m} + \dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz} = \frac{\dot{m}^2}{A^2} v \frac{dv}{dz}$	$-\frac{\dot{m}^2}{A^3} v \frac{dA}{dz} = - \left[\left(\frac{\dot{m}}{A} \right)^2 - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $\dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz} = \frac{\dot{m}^2}{A^2} v \frac{dv}{dz}$
	Constant	$0 = - \left[\left(\frac{\dot{m}}{A} \right)^2 - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $\frac{\dot{q}}{m} = \frac{\dot{m}^2}{A^2} v \frac{dv}{dz}$	$0 = - \left[\left(\frac{\dot{m}}{A} \right)^2 - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $\frac{\dot{q}}{m} = \frac{\dot{m}^2}{A^2} v \frac{dv}{dz}$
Isochor	Variable	$-\frac{\dot{m}^2}{A^3} v \frac{dA}{dz} = -\frac{R_t}{v} \frac{dT}{dz}$ $\frac{\dot{q}}{m} + \dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz} = \frac{\kappa}{\kappa-1} R_t \frac{dT}{dz}$	$-\frac{\dot{m}^2}{A^3} v \frac{dA}{dz} = -\frac{R_t}{v} \frac{dT}{dz}$ $\dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz} = \frac{\kappa}{\kappa-1} R_t \frac{dT}{dz}$
	Constant	$0 = -\frac{R_t}{v} \frac{dT}{dz}$ $\frac{\dot{q}}{m} = \frac{\kappa}{\kappa-1} R_t \frac{dT}{dz}$	$0 = -\frac{R_t}{v} \frac{dT}{dz}$ $0 = \frac{\kappa}{\kappa-1} R_t \frac{dT}{dz}$
Isobar	Variable	$-\frac{\dot{m}^2}{A^3} v \frac{dA}{dz} = - \left(\frac{\dot{m}}{A} \right)^2 \frac{dv}{dz}$ $\frac{\dot{q}}{m} + \dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz} = \left[\frac{\kappa}{\kappa-1} p + \frac{\dot{m}^2}{A^2} v \right] \frac{dv}{dz}$	$-\frac{\dot{m}^2}{A^3} v \frac{dA}{dz} = - \left(\frac{\dot{m}}{A} \right)^2 \frac{dv}{dz}$ $\dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz} = \left[\frac{\kappa}{\kappa-1} p + \frac{\dot{m}^2}{A^2} v \right] \frac{dv}{dz}$
	Constant	$0 = - \left(\frac{\dot{m}}{A} \right)^2 \frac{dv}{dz}$ $\frac{\dot{q}}{m} = \left[\frac{\kappa}{\kappa-1} p + \frac{\dot{m}^2}{A^2} v \right] \frac{dv}{dz}$	$0 = - \left(\frac{\dot{m}}{A} \right)^2 \frac{dv}{dz}$ $0 = \left[\frac{\kappa}{\kappa-1} p + \frac{\dot{m}^2}{A^2} v \right] \frac{dv}{dz}$

Table 2. Wet Steam Flow Equations I.

Table 2. Wet Steam Flow Equations II.

Table 3. Superheated Steam Flow Equations I.

Table 3. Superheated Steam Flow Equations II.

	Cross-Section	No Friction	No Heat Loss No Friction
Polytropic	Variable	$-\frac{\dot{m}^2}{A^3} v \frac{dA}{dz} = -\frac{R_t}{v} \frac{dT}{dz} - \left[\left(\frac{\dot{m}}{A} \right)^2 - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $\dot{q}_m + \dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz} = \left(\frac{\partial h}{\partial T} \right)_v \frac{dT}{dz} + \left[\left(\frac{\partial h}{\partial v} \right)_T - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$	
	Constant	$0 = -\frac{R_t}{v} \frac{dT}{dz} - \left[\left(\frac{\dot{m}}{A} \right)^2 - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $\dot{q}_m = \left(\frac{\partial h}{\partial T} \right)_v \frac{dT}{dz} + \left[\left(\frac{\partial h}{\partial v} \right)_T - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$	
Adiabatic	Variable		$-\frac{\dot{m}^2}{A^3} v \frac{dA}{dz} = -\frac{R_t}{v} \frac{dT}{dz} - \left[\left(\frac{\dot{m}}{A} \right)^2 - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $\dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz} = \left(\frac{\partial h}{\partial T} \right)_v \frac{dT}{dz} + \left[\left(\frac{\partial h}{\partial v} \right)_T + \frac{\dot{m}^2}{A^2} v \right] \frac{dv}{dz}$
	Constant		$0 = -\frac{R_t}{v} \frac{dT}{dz} - \left[\left(\frac{\dot{m}}{A} \right)^2 - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $0 = \left(\frac{\partial h}{\partial T} \right)_v \frac{dT}{dz} + \left[\left(\frac{\partial h}{\partial v} \right)_T - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$
Isentropic	Variable	$-\frac{\dot{m}^2}{A^3} v \frac{dA}{dz} = -\left[-\frac{R_t}{v} \frac{dT}{dz} - \left(\frac{\dot{m}}{A} \right)^2 + \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $\dot{q}_m + \dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz} = \left[\left(\frac{\partial h}{\partial v} \right)_S + \frac{\dot{m}^2}{A^2} v \right] \frac{dv}{dz}$	$-\frac{\dot{m}^2}{A^3} v \frac{dA}{dz} = -\left[-\frac{R_t}{v} \frac{dT}{dz} - \left(\frac{\dot{m}}{A} \right)^2 + \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $\dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz} = \left[\left(\frac{\partial h}{\partial v} \right)_S + \frac{\dot{m}^2}{A^2} v \right] \frac{dv}{dz}$
	Constant	$0 = \left[-\frac{R_t}{v} \frac{dT}{dz} - \left(\frac{\dot{m}}{A} \right)^2 + \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $\dot{q}_m = \left[\left(\frac{\partial h}{\partial v} \right)_S + \frac{\dot{m}^2}{A^2} v \right] \frac{dv}{dz}$	$0 = \left[-\frac{R_t}{v} \frac{dT}{dz} - \left(\frac{\dot{m}}{A} \right)^2 + \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $0 = \left[\left(\frac{\partial h}{\partial v} \right)_S + \frac{\dot{m}^2}{A^2} v \right] \frac{dv}{dz}$
Isothermal	Variable	$-\frac{\dot{m}^2}{A^3} v \frac{dA}{dz} = -\left[\left(\frac{\dot{m}}{A} \right)^2 - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $\dot{q}_m + \dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz} = \left[\left(\frac{\partial h}{\partial v} \right)_T - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$	$-\frac{\dot{m}^2}{A^3} v \frac{dA}{dz} = -\left[\left(\frac{\dot{m}}{A} \right)^2 - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $\dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz} = \left[\left(\frac{\partial h}{\partial v} \right)_T - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$
	Constant	$0 = -\left[\left(\frac{\dot{m}}{A} \right)^2 - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $\dot{q}_m = \left[\left(\frac{\partial h}{\partial v} \right)_T - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$	$0 = -\left[\left(\frac{\dot{m}}{A} \right)^2 - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$ $0 = \left[\left(\frac{\partial h}{\partial v} \right)_T - \frac{R_t T}{v^2} \right] \frac{dv}{dz}$
Isochor	Variable	$-\frac{\dot{m}^2}{A^3} v \frac{dA}{dz} = -\frac{R_t}{v} \frac{dT}{dz}$ $\dot{q}_m + \dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz} = \left(\frac{\partial h}{\partial T} \right)_v \frac{dT}{dz}$	$-\frac{\dot{m}^2}{A^3} v \frac{dA}{dz} = -\frac{R_t}{v} \frac{dT}{dz}$ $\dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz} = \left(\frac{\partial h}{\partial T} \right)_v \frac{dT}{dz}$
	Constant	$0 = -\frac{R_t}{v} \frac{dT}{dz}$ $\dot{q}_m = \left(\frac{\partial h}{\partial T} \right)_v \frac{dT}{dz}$	$0 = -\frac{R_t}{v} \frac{dT}{dz}$ $0 = \left(\frac{\partial h}{\partial T} \right)_v \frac{dT}{dz}$
Isobar	Variable	$-\frac{\dot{m}^2}{A^3} v \frac{dA}{dz} = -\left(\frac{\dot{m}}{A} \right)^2 \frac{dv}{dz}$ $\dot{q}_m + \dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz} = \left[\left(\frac{\partial h}{\partial v} \right)_p + \frac{\dot{m}^2}{A^2} v \right] \frac{dv}{dz}$	$-\frac{\dot{m}^2}{A^3} v \frac{dA}{dz} = -\left(\frac{\dot{m}}{A} \right)^2 \frac{dv}{dz}$ $\dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz} = \left[\left(\frac{\partial h}{\partial v} \right)_p + \frac{\dot{m}^2}{A^2} v \right] \frac{dv}{dz}$
	Variable	$0 = -\left(\frac{\dot{m}}{A} \right)^2 \frac{dv}{dz}$ $\dot{q}_m = \left[\left(\frac{\partial h}{\partial v} \right)_p + \frac{\dot{m}^2}{A^2} v \right] \frac{dv}{dz}$	$0 = -\left(\frac{\dot{m}}{A} \right)^2 \frac{dv}{dz}$ $0 = \left[\left(\frac{\partial h}{\partial v} \right)_p + \frac{\dot{m}^2}{A^2} v \right] \frac{dv}{dz}$

Table 4. Solution of the System od Differential Equations for Ideal Gas with Heat Loss and Friction I.

Cross-Section	Prepared Equations for Solution	Solutions
Polytropic	$a_{11}(z, v) \frac{dT}{dz} + a_{12}(z, v, T) \frac{dv}{dz} = b_{11}(z, v)$ $a_{21} \frac{dT}{dz} + a_{22}(z, v) \frac{dv}{dz} = b_{21}(z, v, T)$ $a_{11}(z, v) = -\frac{\dot{R}_i}{v}, a_{12}(z, v, T) = \left[\left(\frac{\dot{m}}{A} \right)^2 - \frac{R_i T}{v^2} \right]$ $a_{21} = \frac{\kappa-1}{\kappa-1} R_i, a_{22}(z, v) = \left(\frac{\dot{m}}{A} \right)^2 v$ $b_{11}(z, v) = \frac{\lambda}{2d} v \left(\frac{\dot{m}}{A} \right)^2 - \frac{\dot{m}^2}{A^3} v \frac{dA}{dz}$ $b_{21}(z, v, T) = \frac{\dot{q}}{m} + \dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz}$	$\begin{bmatrix} \frac{dT}{dz} \\ \frac{dv}{dz} \end{bmatrix}_z = \begin{bmatrix} a_{11}(z, v) & a_{12}(z, v, T) \\ a_{21} & a_{22}(z, v) \end{bmatrix}_z^{-1} \begin{bmatrix} b_{11}(z, v) \\ b_{21}(z, v, T) \end{bmatrix}_z$ $\begin{bmatrix} T(z + \Delta z) \\ v(z + \Delta z) \end{bmatrix} = \begin{bmatrix} \frac{dT}{dz} \\ \frac{dv}{dz} \end{bmatrix} \Delta z \begin{bmatrix} T(z) \\ v(z) \end{bmatrix}$
Constant	$a_{11}(z, v) \frac{dT}{dz} + a_{12}(z, v, T) \frac{dv}{dz} = b_{11}(z, v)$ $a_{21} \frac{dT}{dz} + a_{22}(z, v) \frac{dv}{dz} = b_{21}(z, v, T)$ $a_{11}(z, v) = -\frac{\dot{R}_i}{v}, a_{12}(z, v, T) = \left[\left(\frac{\dot{m}}{A} \right)^2 - \frac{R_i T}{v^2} \right]$ $a_{21} = \frac{\kappa-1}{\kappa-1} R_i, a_{22}(z, v) = \left(\frac{\dot{m}}{A} \right)^2 v$ $b_{11}(z, v) = \frac{\lambda}{2d} v \left(\frac{\dot{m}}{A} \right)^2$ $b_{21}(z, v, T) = \frac{\dot{q}}{m}$	$\begin{bmatrix} \frac{dT}{dz} \\ \frac{dv}{dz} \end{bmatrix}_z = \begin{bmatrix} a_{11}(z, v) & a_{12}(z, v, T) \\ a_{21} & a_{22}(z, v) \end{bmatrix}_z^{-1} \begin{bmatrix} b_{11}(z, v) \\ b_{21}(z, v, T) \end{bmatrix}_z$ $\begin{bmatrix} T(z + \Delta z) \\ v(z + \Delta z) \end{bmatrix} = \begin{bmatrix} \frac{dT}{dz} \\ \frac{dv}{dz} \end{bmatrix} \Delta z \begin{bmatrix} T(z) \\ v(z) \end{bmatrix}$
	$\frac{dv}{dz} = \frac{-\frac{\dot{q} A^2}{m^3} \frac{\kappa-1}{\kappa} v - \frac{\lambda}{2d} v^3}{\left(1 + \frac{1}{2} \frac{\kappa-1}{\kappa}\right)v^2 - \frac{\dot{q} A^2}{m^3} \frac{\kappa-1}{\kappa} z - \frac{R_i A^2}{m^2} T_0 - \frac{1}{2} \frac{\kappa-1}{\kappa} v_0^2}$ <p>If $\dot{q} = \text{const.}$</p>	$v(z + \Delta z) = \frac{dv}{dz} \Delta z + v(z)$ $T(z + \Delta z) = T(z) + \frac{\dot{q}}{m} \frac{\kappa-1}{\kappa} \frac{\Delta z}{R_i} + \frac{\kappa-1}{\kappa} \frac{1}{R_i} \left(\frac{\dot{m}}{A} \right)^2 [v(z)^2 - v(z + \Delta z)^2]$

Table 4. Solution of the System od Differential Equations for Ideal Gas with Heat Loss and Friction II.

	Cross-Section	Prepared Equations for Solution	Solutions
Adiabatic	Variable	$a_{11}(z, v) \frac{dT}{dz} + a_{12}(z, v, T) \frac{dv}{dz} = b_{11}(z, v)$ $a_{21} \frac{dT}{dz} + a_{22}(z, v) \frac{dv}{dz} = b_{21}(z, v, T)$ $a_{11}(z, v) = -\frac{R_i}{v}, a_{12}(z, v, T) = \left(\frac{\dot{m}}{A}\right)^2 - \frac{R_i T}{v^2}$ $a_{21} = \frac{\kappa}{\kappa-1} R_i, a_{22}(z, v) = \left(\frac{\dot{m}}{A}\right)^2 v$ $b_{11}(z, v) = \frac{\lambda}{2d} v \left(\frac{\dot{m}}{A}\right)^2 - \frac{\dot{m}^2}{A^3} v \frac{dA}{dz}$ $b_{21}(z, v, T) = \dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz}$	$\begin{bmatrix} \frac{dT}{dz} \\ \frac{dv}{dz} \end{bmatrix}_z = \begin{bmatrix} a_{11}(z, v) & a_{12}(z, v, T) \\ a_{21} & a_{22}(z, v) \end{bmatrix}_z^{-1} \begin{bmatrix} b_{11}(z, v) \\ b_{21}(z, v, T) \end{bmatrix}_z$ $\begin{bmatrix} T(z + \Delta z) \\ v(z + \Delta z) \end{bmatrix} = \begin{bmatrix} \frac{dT}{dz} \\ \frac{dv}{dz} \end{bmatrix}_{\Delta z} \begin{bmatrix} T(z) \\ v(z) \end{bmatrix}$
	Constant	$a_{11}(z, v) \frac{dT}{dz} + a_{12}(z, v, T) \frac{dv}{dz} = b_{11}(z, v)$ $a_{21} \frac{dT}{dz} + a_{22}(z, v) \frac{dv}{dz} = b_{21}(z, v, T)$ $a_{11}(z, v) = -\frac{R_i}{v}, a_{12}(z, v, T) = \left(\frac{\dot{m}}{A}\right)^2 - \frac{R_i T}{v^2}$ $a_{21} = \frac{\kappa}{\kappa-1} R_i, a_{22}(z, v) = \left(\frac{\dot{m}}{A}\right)^2 v$ $b_{11}(z, v) = \frac{\lambda}{2d} v \left(\frac{\dot{m}}{A}\right)^2$ $b_{21}(z, v, T) = 0$	$\begin{bmatrix} \frac{dT}{dz} \\ \frac{dv}{dz} \end{bmatrix}_z = \begin{bmatrix} a_{11}(z, v) & a_{12}(z, v, T) \\ a_{21} & a_{22}(z, v) \end{bmatrix}_z^{-1} \begin{bmatrix} b_{11}(z, v) \\ b_{21}(z, v, T) \end{bmatrix}_z$ $\begin{bmatrix} T(z + \Delta z) \\ v(z + \Delta z) \end{bmatrix} = \begin{bmatrix} \frac{dT}{dz} \\ \frac{dv}{dz} \end{bmatrix}_{\Delta z} \begin{bmatrix} T(z) \\ v(z) \end{bmatrix}$
		$\frac{dv}{dz} = \frac{-\frac{\lambda}{2d} v^3}{\left(1 + \frac{1}{2} \frac{\kappa-1}{\kappa}\right) v^2 - \frac{R_i A^2}{\dot{m}^2} T_0 - \frac{1}{2} \frac{\kappa-1}{\kappa} v_0^2}$ <p>If $\dot{q} = \text{const.}$</p>	$v(z + \Delta z) = \frac{dv}{dz} \Delta z + v(z)$ $T(z + \Delta z) = T(z) + \frac{\kappa-1}{\kappa} \frac{1}{R_i} \left(\frac{\dot{m}}{A}\right)^2 [v(z)^2 - v(z + \Delta z)^2]$

Table 4. Solution of the System of Differential Equations for Ideal Gas with Heat Loss and Friction III.

	Cross-Section	Prepared Equations for Solution	Solutions
Isentropic Variable		$\frac{dv}{dz} = \frac{\left(\frac{\lambda}{2d} - \frac{1}{A} \frac{dA}{dz}\right)v^{\kappa+2}}{\frac{\kappa R_i T_0 v_0^{\kappa-1}}{\dot{m}^2} A^2 - v^{\kappa+1}}$ $\frac{dv}{dz} = \frac{\frac{\kappa-1}{R_i K} \frac{\dot{q}}{m} v^\kappa + \frac{\kappa-1}{R_i K} \dot{m} \frac{1}{A^3} \frac{dA}{dz} v^{\kappa+2}}{(1-\kappa) T_0 v_0^{\kappa-1} + \left(\frac{\dot{m}}{A}\right)^2 \frac{\kappa-1}{R_i K} v^{\kappa+1}}$	<p>If $A(z)$ is given, the heat loss $\dot{q}(z) = -\left(\frac{\lambda}{2d} - \frac{1}{A} \frac{dA}{dz}\right) \left(\frac{\dot{m}}{A}\right)^2 v^2 \dot{m}$ and $v(z)$</p> $\frac{dv}{dz} = \frac{\left(\frac{\lambda}{2d} - \frac{1}{A} \frac{dA}{dz}\right)v^{\kappa+2}}{\frac{\kappa R_i T_0 v_0^{\kappa-1}}{\dot{m}^2} A^2 - v^{\kappa+1}}$ <p>If $\dot{q}(z)$ is given</p> $\left(\frac{\kappa R_i T_0 v_0^{\kappa-1}}{\dot{m}^2} A^2 - v^{\kappa+1}\right) \frac{dv}{dz} + \frac{1}{A} v^{\kappa+2} \frac{dA}{dz} = \frac{\lambda \sqrt{\pi}}{4\sqrt{A}} v^{\kappa+2}$ $\left[\left(\frac{\dot{m}}{A}\right)^2 v^{\kappa+1} - \kappa R_i T_0 v_0^{\kappa-1}\right] \frac{dv}{dz} - \frac{\dot{m}}{A^3} v^{\kappa+2} \frac{dA}{dz} = v^\kappa \frac{\dot{q}}{m}$
	Constant	$\frac{dv}{dz} = \frac{\frac{\lambda}{2d} v^{\kappa+2}}{\frac{\kappa R_i T_0 v_0^{\kappa-1}}{\dot{m}^2} A^2 - v^{\kappa+1}}$ $\frac{dv}{dz} = \frac{\frac{\kappa-1}{R_i K} \frac{\dot{q}}{m} v^\kappa}{(1-\kappa) T_0 v_0^{\kappa-1} + \left(\frac{\dot{m}}{A}\right)^2 \frac{\kappa-1}{R_i K} v^{\kappa+1}}$	<p>The heat loss</p> $\dot{q}(z) = -\frac{\lambda}{2d} \left(\frac{\dot{m}}{A}\right)^2 v^2 \dot{m}^2$ and $v(z)$ $\frac{R_i K T_0 v_0^{\kappa-1}}{\dot{m}^{2(\kappa+1)} A^2} A^2 \left(\frac{1}{v_0^{\kappa+2}} - \frac{1}{v^{\kappa+2}}\right) - \ln \frac{V_0}{V} = \frac{\lambda}{2d} z$

Table 4. Solution of the System of Differential Equations for Ideal Gas with Heat Loss and Friction IV.

	Cross-Section	Prepared Equations for Solution	Solution
Isothermal	Variable	$\frac{dv}{dz} = \frac{\left(\frac{\lambda}{2d} - \frac{1}{A} \frac{dA}{dz}\right)v^3}{\frac{R_i T}{\left(\frac{\dot{m}}{A}\right)^2} - v^2}$ $\frac{dv}{dz} = \frac{\dot{q}}{\dot{m}^3} A^2 \frac{1}{V} + v \frac{1}{A} \frac{dA}{dz}$	<p>If $A(z)$ is given, the \dot{q} is</p> $\dot{q}(z) = \left[\frac{\left(\frac{\lambda}{2d} - \frac{1}{A} \frac{dA}{dz}\right)v^3}{\frac{R_i T}{\left(\frac{\dot{m}}{A}\right)^2} - v^2} - v \frac{1}{A} \frac{dA}{dz} \right] \frac{\dot{m}^3}{A^2} v$ <p>for $v(z)$</p> $\frac{dv}{dz} = \frac{\left(\frac{\lambda\sqrt{\pi}}{4\sqrt{A}} - \frac{1}{A} \frac{dA}{dz}\right)v^3}{\frac{R_i T}{\left(\frac{\dot{m}}{A}\right)^2} A^2 - v^2}$ <p>If $\dot{q}(z)$ is given</p> $\left(\frac{R_i T A^2}{\dot{m}^2} - v^2\right) \frac{dv}{dz} + \frac{1}{A} v^3 \frac{dA}{dz} = \frac{\lambda\sqrt{\pi}}{4\sqrt{A}} v^3$ $\frac{dv}{dz} - \frac{v}{A} \frac{dA}{dz} = \frac{\dot{q}}{\dot{m}^3} A^2 \frac{1}{V}$
	Constant	$\frac{dv}{dz} = \frac{\frac{\lambda}{2d} v^3}{\frac{R_i T}{\left(\frac{\dot{m}}{A}\right)^2} - v^2}$ $\frac{dv}{dz} = \frac{\dot{q}}{\dot{m}^3} A^2 \frac{1}{V}$	<p>If $A(z)$ is given, the \dot{q} is</p> $\dot{q}(z) = \frac{\frac{\lambda}{2d} v^4}{\frac{R_i T}{\left(\frac{\dot{m}}{A}\right)^2} - v^2} \frac{\dot{m}^3}{A^2}$ <p>for $v(z)_-$</p> $\frac{1}{2} \frac{R_i T}{\left(\frac{\dot{m}}{A}\right)^2} \left(\frac{1}{v_0^2} - \frac{1}{v^2} \right) + \ln \frac{v_0}{v} = \frac{\lambda}{2d} z$
Isochor	Variable	$\frac{dT}{dz} = - \left[\frac{\lambda}{2d} \left(\frac{\dot{m}}{A} \right)^2 - \frac{\dot{m}^2}{A^3} \frac{dA}{dz} \right] \frac{1}{R_i} v^2$ $\frac{dT}{dz} = \frac{\dot{q}}{\dot{m}^3} A^2 \frac{1}{V} + v \frac{1}{A} \frac{dA}{dz}$	<p>If $A(z)$ is given, the \dot{q} is</p> $\dot{q}(z) = \frac{\dot{m}^3}{A^3} v^2 \frac{dA}{dz} \frac{1}{\kappa-1} - \frac{\lambda}{2d} \frac{\dot{m}^3}{A^2} \frac{\kappa}{\kappa-1} v^2$ <p>If $\dot{q}(z)$ is given</p> $\frac{dA}{dz} = \frac{\dot{q}(z)(\kappa-1)}{\dot{m}^3 v^2} A^3 + \frac{\lambda\sqrt{\pi}}{4} \kappa v^2 \sqrt{A}$
	Constant	$\frac{dT}{dz} = - \frac{\lambda}{2d} v^2 \left(\frac{\dot{m}}{A} \right)^2 \frac{1}{R_i}$ $\frac{dT}{dz} = \frac{\dot{q}}{m \kappa R_i}$	<p>The heat loss</p> $\dot{q}(z) = - \frac{\lambda}{2d} \frac{\dot{m}^3}{A^2} \frac{\kappa}{\kappa-1} v^2$

Table 4. Solution of the System of Differential Equations for Ideal Gas with Heat Loss and Friction V.

	Cross-Section	Prepared Equations for Solution	Solution
Isobar	Variable	$\frac{dv}{dz} = \left(\frac{1}{A} \frac{dA}{dz} - \frac{\lambda}{2d} \right) v$ $\frac{dv}{dz} = \frac{\frac{\dot{q}}{m} + \dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz}}{\frac{\kappa-1}{\kappa} p + \left(\frac{\dot{m}}{A} \right)^2 v}$	<p>If $A(z)$ is given, the heat loss</p> $\frac{\dot{q}}{m} = \left(\frac{\kappa-1}{\kappa} p + \left(\frac{\dot{m}}{A} \right)^2 v \right) \frac{dv}{dz} - \dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz}$ $v = v_0 e^{\frac{1}{A} \frac{dA}{dz} - \frac{\lambda}{2d} z}$ <p>If $\dot{q}(z)$ is given, for $A(z)$ and $v(z)$</p> $\left(\frac{\dot{m}}{A} \right)^2 \frac{dv}{dz} - \frac{\dot{m}^2}{A^3} v \frac{dA}{dz} = \frac{\lambda}{2d} v \left(\frac{\dot{m}}{A} \right)^2$ $\left(\frac{\kappa-1}{\kappa} p \frac{\dot{m}}{A^2} v \right) \frac{dv}{dz} - \dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz} = \frac{\dot{q}}{m}$
	Constant	$\frac{dv}{dz} = -\frac{\lambda}{2d} v$ $\frac{dv}{dz} = \frac{\dot{q}}{m}$ $\frac{dv}{dz} = \frac{\kappa-1}{\kappa} p + \left(\frac{\dot{m}}{A} \right)^2 v$	$v = v_0 e^{-\frac{\lambda}{2d} z}$ $\dot{q} = -\frac{\lambda}{2d} \dot{m} \left[\frac{\kappa-1}{\kappa} p + \left(\frac{\dot{m}}{A} \right)^2 v_0 e^{-\frac{\lambda}{2d} z} \right] v_0 e^{-\frac{\lambda}{2d} z}$

Table 5. Solution of the System of Differential Equations for Wet Steam with Friction and Heat Loss I.

	Cross-Section	Prepared Equations for Solution	Solutions
Polytropic	Variable	$\begin{bmatrix} \frac{dT}{dz} \\ \frac{dv}{dz} \end{bmatrix} = \begin{bmatrix} a_{11}(T) & a_{12}(z) \\ a_{21}(T) & a_{22}(z, v, T) \end{bmatrix} \begin{bmatrix} b_{11}(z, v) \\ b_{21}(z, v, T) \end{bmatrix}$	$\begin{bmatrix} T(z + \Delta z) \\ v(z + \Delta z) \end{bmatrix} = \begin{bmatrix} \frac{dT}{dz} \\ \frac{dv}{dz} \end{bmatrix} \Delta z + \begin{bmatrix} T(z) \\ v(z) \end{bmatrix}$
	Constant	$\begin{bmatrix} \frac{dT}{dz} \\ \frac{dv}{dz} \end{bmatrix} = \begin{bmatrix} a_{11}(T) & a_{12}(z) \\ a_{21}(T) & a_{22}(z, v, T) \end{bmatrix} \begin{bmatrix} b_{11}(z, v) \\ b_{21}(z, v, T) \end{bmatrix}$	$\begin{bmatrix} T(z + \Delta z) \\ v(z + \Delta z) \end{bmatrix} = \begin{bmatrix} \frac{dT}{dz} \\ \frac{dv}{dz} \end{bmatrix} \Delta z + \begin{bmatrix} T(z) \\ v(z) \end{bmatrix}$
		$\frac{dv}{dz} = \frac{\frac{\lambda}{2d}v - \frac{\dot{q}}{m^2}A^2\frac{\delta_1}{b_1}}{\frac{\delta_1}{b_1}v + \frac{A^2}{m^2}\frac{\delta_1}{b_1}g_1 - 1}$	$v = v_0 + \frac{C_1}{B_0}z + \left(\frac{A_0}{C_0} + \frac{B_1}{B_0}\right) \ln \frac{C_1 v_0 - A_0}{C_1 r - A_0}$
Adiabatic	Variable		
Isentropic	Variable	$\frac{dv}{dz} = \frac{\frac{\lambda}{2d}\left(\frac{\dot{m}}{A}\right)^2 - \frac{\dot{m}^2}{A^3}\frac{dA}{dz}v}{-(\mu_1 + 2\mu_2 v + \dots) - \left(\frac{\dot{m}}{A}\right)^2}$ $\frac{dv}{dz} = \frac{\frac{\dot{q}}{m} + \dot{m}v^2\frac{1}{A^3}\frac{dA}{dz}}{(\gamma_1 + 2\gamma_2 v + \dots) + \left(\frac{\dot{m}}{A}\right)^2 v}$	$\dot{q}(z) = \frac{\frac{\dot{m}^2}{A^3}\frac{dA}{dz} - \frac{\lambda}{2d}\left(\frac{\dot{m}}{A}\right)^2}{(\mu_1 + 2\mu_2 v + \dots) + \left(\frac{\dot{m}}{A}\right)^2} \dot{m} - \frac{\dot{m}^3 v^2}{A^3} \frac{dA}{dz}, \quad \text{If } A = A(z) \text{ is given.}$ $\mu_1 \ln \frac{v}{v_0} + \left[2\mu_2 + \left(\frac{\dot{m}}{A}\right)^2\right] (v - v_0) + \frac{3\mu_3}{2} (v^2 - v_0^2) + \dots = \int_0^z \left[\frac{\dot{m}^2}{A^3} \frac{dA}{dz} - \frac{\lambda}{2d} \left(\frac{\dot{m}}{A}\right)^2 \right] dz$ $\frac{dA}{dz} = \left[\frac{\dot{q} \left[\mu_1 + 2\mu_2 v + \dots + \left(\frac{\dot{m}}{A}\right)^2 \right]}{[\gamma_1 + 2\gamma_2 v + \dots + \left(\frac{\dot{m}}{A}\right)^2] \dot{m}} + \frac{\lambda}{2d} \left(\frac{\dot{m}}{A}\right)^2 \right] \frac{A^3}{\dot{m}^2}$ $A(z + \Delta z) = \frac{dA}{dz} \Delta z + A(z), \quad \dot{q} = \dot{q}(z) \text{ is given}$
	Constant	$\frac{dv}{dz} = \frac{\frac{\lambda}{2d}\left(\frac{\dot{m}}{A}\right)^2 v}{-(\mu_1 + 2\mu_2 v + \dots) - \left(\frac{\dot{m}}{A}\right)^2}$ $\frac{dv}{dz} = \frac{\frac{\dot{q}}{m}}{(\mu_1 + 2\mu_2 v + \dots) + \left(\frac{\dot{m}}{A}\right)^2 v}$	$\dot{q}(z) = -\frac{\lambda}{2d} \left(\frac{\dot{m}}{A}\right)^2 v \frac{v[(\gamma_1 + 2\gamma_2 v + \dots) - \left(\frac{\dot{m}}{A}\right)^2]}{(\mu_1 + 2\mu_2 v + \dots) + \left(\frac{\dot{m}}{A}\right)^2}$ $- \frac{\lambda}{2d} \left(\frac{\dot{m}}{A}\right)^2 \frac{1}{\mu_1 - \left(\frac{\dot{m}}{A}\right)^2} z, \quad \text{If } \mu_2 = \mu_3 = \dots = 0$ $v = v_0$

Table 5. Solution of the System of Differential Equations for Wet Steam with Friction and Heat Loss II.

Cross-S.		Comments and Notations		
Polytropic	Variable	$a_{11}(T) \approx -(\delta_1 + 2\delta_2 T + \dots)$ $a_{12}(z) = -\left(\frac{\dot{m}}{A}\right)^2$	$a_{21}(T, v) = \left(\frac{\partial h}{\partial T}\right)_V$ $a_{22}(z, T, v) = \left[\left(\frac{\partial h}{\partial v}\right)_T + \left(\frac{\dot{m}}{A}\right)^2 v\right]$	$b_{11} = \frac{\lambda}{2d} v \left(\frac{\dot{m}}{A}\right)^2 - \frac{\dot{m}^2}{A^3} v \frac{dA}{dz}$ $b_{21}(z, v, T) = \frac{\dot{q}}{m} + \dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz}$ $h(v = \text{const.}, T) = b_0 + b_1 T + \dots$ $h(T = \text{const.}) = g_0 + g_1 v + \dots$ $p = \delta_0 + \delta_1 T + \dots$
	Constant	$\frac{dA}{dz} = 0$	$a_{11}(T) \approx -(\delta_1 + 2\delta_2 T + \dots)$ $a_{12}(z) = -\left(\frac{\dot{m}}{A}\right)^2$	$a_{21}(T, v) = \left(\frac{\partial h}{\partial T}\right)_V$ $a_{22}(z, T, v) = \left[\left(\frac{\partial h}{\partial v}\right)_T + \left(\frac{\dot{m}}{A}\right)^2 v\right]$ $b_{11}(z, v) = \frac{\lambda}{2d} v \left(\frac{\dot{m}}{A}\right)^2$ $b_{21}(z, v) = \frac{\dot{q}}{m}$
			$\frac{dA}{dz} = 0, \quad \dot{q} = \text{const.}, \quad A_0 = \frac{\dot{q}}{\dot{m}^3} \frac{\delta_1}{b_1} A^2, \quad B_0 = \frac{\delta_1}{b_1}, \quad B_1 = \frac{g_1}{b_1} \delta_1 \frac{A^2}{\dot{m}^2} - 1, \quad C_1 = \frac{\lambda}{2d}, \quad p \approx \delta_0 + \delta_1 T$	
Adiabatic	Variable			
Isentropic	Variable			In the motion equation $\frac{dp}{dz} = \left(\frac{\partial p}{\partial S}\right)_v \frac{dS}{dz} + \left(\frac{\partial p}{\partial v}\right)_S \frac{dv}{dz} = \left(\frac{\partial p}{\partial v}\right)_S \frac{dv}{dz}$ $p \approx \delta_0 + \delta_1 + \dots$ $p(v, S = \text{const.}) \approx \mu_0 + \mu_1 v + \mu_2 v^2 + \dots$ $T(v, S = \text{const.}) \approx v_0 + v_1 v + v_2 v^2 + \dots$ $h(v, S = \text{const.}) \approx \gamma_0 + \gamma_1 v + \gamma_2 v^2 + \dots$
	Constant			If $\mu_2 \equiv \mu_3 \equiv \dots \equiv 0$.

Table 5. Solution of the System of Differential Equations for Wet Steam with Friction and Heat Loss III.

	Cross-Section	Prepared Equations for Solution	Solutions
Isothermal	Variable	$\frac{dv}{dz} = -\left(\frac{\lambda}{2d} - \frac{1}{A} \frac{dA}{dz}\right)v$ $\frac{dv}{dz} = \frac{\frac{\dot{q}}{m} + m^2 \frac{1}{A^3} \frac{dA}{dz} v^2}{(\gamma_1 + 2\gamma_2 v + \dots) + \left(\frac{\dot{m}}{A}\right)^2 v}$	If $A = A(z)$ is given, $\dot{q}(z) = \dot{m} \left(\frac{1}{A} \frac{dA}{dz} - \frac{\lambda}{2d} \right) (\gamma_1 + 2\gamma_2 v + \dots) + \dots$ If $\mu_2 \equiv \mu_3 \equiv \dots \equiv 0$, $v(z) = v_0 \frac{A}{A_0} e^{-\frac{\lambda}{2d} z}$
	Constant	$\frac{dv}{dz} = -\frac{\lambda}{2d} v$ $\frac{dv}{dz} = \frac{\frac{\dot{q}}{m}}{(\gamma_1 + 2\gamma_2 v + \dots) + \left(\frac{\dot{m}}{A}\right)^2 v}$	$\dot{q}(z) = -\frac{\lambda}{2d} v \dot{m} \left[(\gamma_1 + 2\gamma_2 v + \dots) + \left(\frac{\dot{m}}{A}\right)^2 v \right]$ $v(z) = v_0 e^{-\frac{\lambda}{2d} z}$
Isochor	Variable	$\frac{dT}{dz} = \frac{\frac{\lambda}{2d} v \left(\frac{\dot{m}}{A}\right)^2 - \frac{\dot{m}^2}{A^3} v \frac{dA}{dz}}{-(\delta_1 + 2\delta_2 T + \dots)}$ $\frac{dT}{dz} = \left(\frac{\dot{q}}{m} + \dot{m}^2 v^2 \frac{1}{A^3} \frac{dA}{dz} \right) \frac{1}{(b_1 + 2b_2 v + \dots)}$	If $A(z)$ is given, $\dot{q}(z) = \left(\frac{\dot{m}}{A}\right)^3 v \frac{dA}{dz} - \frac{\lambda}{2d} v \left(\frac{\dot{m}^3}{A^2}\right) \frac{b_1 + 2b_2 T + \dots}{\delta_1 + 2\delta_2 T + \dots} - \dot{m}^3 v^2 \frac{1}{A^3} \frac{dA}{dz}$ $T_0 - T + \frac{\delta_2}{\delta_1} (T_0^2 - T^2) + \dots = \frac{1}{\delta_1} \int_0^z \left(\frac{\lambda}{2d} \dot{m}^2 v \frac{1}{A^2} - \dot{m}^2 v \frac{1}{A^3} \frac{dA}{dz} \right) dz$
	Constant	$\frac{dT}{dz} = \frac{\frac{\lambda}{2d} v \left(\frac{\dot{m}}{A}\right)^2}{-(\delta_1 + 2\delta_2 T + \dots)}$ $\frac{\dot{q}}{m} = \frac{1}{(b_1 + 2b_2 v + \dots)}$	$\dot{q}(z) = -\frac{\lambda}{2d} v \left(\frac{\dot{m}^3}{A^2}\right) \frac{b_1 + 2b_2 T + \dots}{\delta_1 + 2\delta_2 T + \dots}$ $T_0 - T + \frac{\delta_2}{\delta_1} (T_0^2 - T^2) + \dots = \frac{1}{\delta_1} \frac{\lambda}{2d} \frac{\dot{m}^2}{A^2} v z$
Isobar	Variable	As Isothermal flow	
	Constant	As Isothermal Flow	

Table 5. Solution of the System of Differential Equations for Wet Steam with Friction and Heat Loss IV.

		Comments and Notations
Isothermal	Variable	$h(T = \text{const.}, v) = g_0 + g_1 T + g_2 T^2 + \dots$
	Constant	$h(T = \text{const.}, v) = g_0 + g_1 T + g_2 T^2 + \dots$ $\frac{dA}{dz} = 0$
Isochor	Variable	$h(v = \text{const.}, T) = b_0 + b_1 T + b_2 T^2 + \dots$ $p = \delta_0 + \delta_1 T + \dots$
	Constant	$h(v = \text{const.}, T) = b_0 + b_1 T + b_2 T^2 + \dots$ $\frac{dA}{dz} = 0$
Isobar	Variable	As Isothermal flow
	Constant	As Isothermal Flow

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