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RESEARCH ARTICLE

Two Theorems for Computation of Projections of Virtual Displacements and Its Application in Structural Analysis

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Abstract

A theorem for planar case and its generalization for spatial case are proposed to determine the projection of a virtual displacement to the orientation under the case of knowing the projections of a virtual displacement to the given two or three orientations for object systems subject to holonomic and scleronomic constraints. Some lemmas corresponding to the two theorems for special cases are given. Applications to structural static analysis are investigated using the two theorems in this paper. Result reveals that the two theorems and corresponding lemmas are easy to be used, shorten the distance between the principle of virtual displacement and its application, and the relating problems can be solved quickly with them.

Keywords

projection of virtual displacement, principle of virtual displacement, engineering mechanics, structural analysis

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1 Introduction It is well known that the

It is well known that the principle of virtual displacement (or virtual work) is a main part in analytical statics, and the important basis for analytical dynamics and structural analysis. This principle provides an excellent tool for people to investigate the equilibrium laws of object systems, and plays a supporting role in classical mechanics.

Generally speaking, some current college textbooks [1-7] for engineering mechanics or structural analysis course often propose two main methods for building relations among virtual displacements of different points for the system with holonomic and scleronomic constraints, ie, analytical method and geometrical method. Using analytical method, one can determine the virtual displacements of interest points (usually points of forces) by taking differentials of their position coordinates, and geometrical methods by introducing the concept of the instant center of rotation of virtual displacements. These methods are very useful in application of the principle of virtual displacement. However, in actual experience, we realize that these methods are sometimes not sufficient for analyzing complex structures. For some difficult problems, such as the examples given in this paper, it is impossible or not sufficient to solve them out if only directly using the results given in these current textbooks. Then, how to shorten the distance between the principle of virtual displacement and its application, and give some feasible approaches to realize, is very important and great significance.

In this study, we will gain an insight into the application of the principle of virtual displacement and present two new theorems on how to build up the relations among the projections to different orientations of the virtual displacement of a point. Several examples are given to illustrate the application of two theorems and corresponding lemmas. The article is organized as follows. In Section 2, a theorem is given to find a projection under knowing the two projections to different orientations of a point for planar systems, and two lemmas for special cases are proposed. In Section 3, a theorem for spatial systems as the planar extending situation is given to find a projection under knowing the three projections to different orientations of a point, and three lemmas for special cases are proposed.

2 The first theorem for plane cases

For a planar object system, if the projections of a virtual displacement δr of a certain dot to orientations n_1 and n_2 are δr_1 and δr_2 separately, shown in Fig. 1, then the projection of the virtual displacement to orientation n_3 is

$$\delta r_3 = \frac{\delta r_1 \sin(\theta - \phi) + \delta r_2 \sin\phi}{\sin\theta},\tag{1}$$

where, angles θ and φ are angles between \mathbf{n}_1 and \mathbf{n}_2 , \mathbf{n}_1 and \mathbf{n}_3 separately.



Fig. 1 A virtual displacement and its projections

Proof. Assume e, e_1 , e_2 and e_3 are the unit vectors of the orientations of n, n_1 , n_2 and n_3 , separately. Take the start point A of the virtual displacement δr as the origin of the coordinate system xAy. Let a is the angle between n_1 and x axis. According to the given conditions, one reads

$$\delta \boldsymbol{r} \cdot \boldsymbol{e}_1 = \delta r_1,$$
$$\delta \boldsymbol{r} \cdot \boldsymbol{e}_2 = \delta r_2.$$

Hence, one can get

$$\delta \boldsymbol{r} \cdot \left(\cos\alpha \boldsymbol{i} + \sin\alpha \boldsymbol{j}\right) = \delta r_1 \tag{2}$$

$$\delta \boldsymbol{r} \cdot \left[\cos\left(\alpha + \theta\right) \boldsymbol{i} + \sin\left(\alpha + \theta\right) \boldsymbol{j} \right] = \delta r_2 \tag{3}$$

where *i*, *j* are the standard unit vectors of *x* and *y* axis.

By taking the inner products $\delta \mathbf{r} \cdot \mathbf{i}$ and $\delta \mathbf{r} \cdot \mathbf{j}$ as unknowns, and solving the equations of (2) and (3), one can get

$$\delta \boldsymbol{r} \cdot \boldsymbol{i} = \frac{\delta r_1 \sin\left(\alpha + \theta\right) - \delta r_2 \sin\alpha}{\sin\theta} \tag{4}$$

$$\delta \mathbf{r} \cdot \mathbf{j} = \frac{\delta r_2 \cos \alpha - \delta r_1 \cos \left(\alpha + \theta\right)}{\sin \theta}$$
(5)

Then, the projection of the virtual displacement to orientation n_3 is

$$\delta r_{3} = \delta \boldsymbol{r} \cdot \boldsymbol{e}_{3} = \delta \boldsymbol{r} \cdot \left[\cos(\alpha + \phi) \boldsymbol{i} + \sin(\alpha + \phi) \boldsymbol{j} \right]$$
$$= \cos(\alpha + \phi) \cdot \frac{\delta r_{1} \sin(\alpha + \phi) - \delta r_{2} \sin\alpha}{\sin\theta}$$
$$+ \sin(\alpha + \phi) \cdot \frac{\delta r_{2} \cos\alpha - \delta r_{1} \cos(\alpha + \phi)}{\sin\theta}$$
$$= \frac{\delta r_{1} \sin(\alpha + \phi) + \delta r_{2} \sin\phi}{\sin\theta}.$$

Lemma 1. When $\theta \neq 0(180^\circ)$, $\delta r_1 = \delta r_2 = 0$, then $\delta r_3 = 0$. **Lemma 2.** When $\theta = 90^\circ$, then $\delta r_3 = \delta r_1 \cos \varphi + \delta r_2 \sin \varphi$.

For demonstrating the effectiveness of this theorem in solving the balancing problems of planar object systems, two examples are given below.

Example 1. Determine the force in member *CD* of the truss, shown in Fig. 2. Assume all members are pin connected, and AD = AE = ED = EC = CG = DG = DB = BG = a.



Fig. 2 Example 1

Solution. To calculate the force in rod *CD*, it is isolated from the system in Fig. 2 (b). Then, the system in Fig. 2 (b) is a mechanism. Obviously, equilateral triangles *AED* and *DBG* can be regarded as rigid plates, and ΔAED can be assumed to rotate about *A* point with $\delta\theta$. Then, $\delta r_E = \delta r_D = a\delta\theta$. Considering the orientation of δr_D and character of support *B*, point *B* is the virtual displacement center of ΔDBG , therefore $\delta r_G = \delta r_D = a\delta\theta$. Based on the principle of virtual displacement, one reads

$$\boldsymbol{F} \cdot \delta \boldsymbol{r}_{G} + \boldsymbol{F}_{CD}' \cdot \delta \boldsymbol{r}_{D} + \boldsymbol{F}_{CD} \cdot \delta \boldsymbol{r}_{C} = 0$$
(6)

Based on the theorem of projection of virtual displacement (i.e. the projections of the virtual displacements of the points from a straight line belonging to a body, on that line, are equal.), the projections of δr_C to the orientations $E \rightarrow C$ and $G \rightarrow C$ are zeros. Then, by employing the lemma 1 of the above theorem, δr_C must be zero. Therefore, based on the above analysis and (6), one can get

$$-F \cdot a\delta\theta \cdot \cos 30^\circ - F'_{CD} \cdot a\delta\theta = 0.$$

Because of $\theta \neq 0$, then

$$F'_{CD} = F_{CD} = -F \cdot \cos 30^\circ = -\frac{\sqrt{3}}{2}F.$$

(bar CD in compression).

Example 2. Determine the force in member DG of the truss, shown in Fig. 3. Assume all members are pin connected.



Fig. 3 Example 2

Solution. To calculate the force in rod *DG*, it is isolated from the system in Fig. 3 (b). Obviously, triangles *ACE* and *DBE* can be regarded as rigid plates, and ΔACE can be assumed to rotate about *A* point with $\delta\theta$. Then, $\delta r_C = 4a\delta\theta$, $\delta r_E = 4\sqrt{2}a\delta\theta$. Considering that the orientation of δr_B is horizontal, the virtual displacement δr_D must be horizontal, and

$$\delta r_D = \delta r_E \cdot \cos 45^\circ = 4a\delta\theta,$$

$$\delta r_B \cdot \cos 45^\circ = \delta r_E, \ \delta r_B = 8a\delta\theta.$$

For the bar BH, it is easy to get that

$$\delta r_H \cdot \cos(90^\circ - 2\alpha) = \delta r_B \cdot \cos\alpha$$

$$\delta r_{H} = \frac{1}{2\sin\alpha} \,\delta r_{B} = \frac{4a}{\sin\alpha} \,\delta\theta$$

Then, for the bar *GH*, the projection of $\delta \mathbf{r}_G$ to the orientation $G \rightarrow H$ is

$$\delta r_{GH} = \delta r_H \cdot \cos \alpha = \frac{4a\cos \alpha}{\sin \alpha} \delta \theta.$$

For the bar CG, the projection of $\delta \mathbf{r}_G$ to the orientation $C \rightarrow G$ is

$$\delta r_{CG} = \delta r_C \cdot \cos(90^\circ - \beta) = 4a \sin \beta \delta \theta$$

For the point *G*, the projections δr_{GH} and δr_{CG} are known. By employing the formula (1) of the above theorem, the projection of δr_G to the orientation $G \rightarrow D$ is

$$\delta r_{GD} = \frac{\delta r_{GH} \sin \left\{ \beta - \left[\beta + 2(90^{\circ} - \beta) \right] \right\} + \delta r_{CG} \sin \left[\beta + 2(90^{\circ} - \beta) \right]}{\sin \beta}$$
$$= \left[\frac{-4a \cos \alpha \sin 2\beta}{\sin \alpha \sin \beta} + 4a \sin \beta \right] \delta \theta$$
$$= \left(-8a \cot \alpha \cos \beta + 4a \sin \beta \right) \delta \theta$$
$$= \left(-8a \cdot 4 \cdot \frac{2}{2\sqrt{5}} + 4a \cdot \frac{4}{2\sqrt{5}} \right) \delta \theta = -\frac{24}{\sqrt{5}} a \delta \theta.$$

By employing the principle of virtual displacement, one can get the coefficient matrix is reversible. Hence, one can get

$$\boldsymbol{F} \cdot \delta \boldsymbol{r}_{D} + \boldsymbol{F}_{DG} \cdot \delta \boldsymbol{r}_{D} + \boldsymbol{F}_{DG}' \cdot \delta \boldsymbol{r}_{G} = 0 \tag{7}$$

Therefore, based on the above analysis and (7), one can get

$$F \cdot 4a\delta\theta - F_{DG} \cdot 4a\delta\theta \cdot \cos(90^\circ - \beta) - F'_{DG} \cdot \frac{24}{\sqrt{5}}a\delta\theta = 0.$$

Because $\delta \theta \neq 0$, $F_{DG} = F'_{DG}$, then

$$F_{DG} = \frac{F}{\sin\beta + \frac{6}{\sqrt{5}}} = \frac{F}{\frac{2}{\sqrt{5}} + \frac{6}{\sqrt{5}}} = \frac{\sqrt{5}F}{8} = 0.279F.$$

3 The second theorem for spatial cases

If the projections of the virtual displacement δr of a dot to three non-coplanar orientations n_1 , n_2 , n_3 are δr_1 , δr_2 , δr_3 , then the projection δr_4 of the virtual displacement to the orientation n_4 is

$$\delta r_{4} = \left(e_{4x}, e_{4y}, e_{4z}\right) \begin{bmatrix} e_{1x} & e_{1y} & e_{1z} \\ e_{2x} & e_{2y} & e_{2z} \\ e_{3x} & e_{3y} & e_{3z} \end{bmatrix}^{-1} \begin{pmatrix} \delta r_{1} \\ \delta r_{2} \\ \delta r_{3} \end{pmatrix}.$$
(8)

where e_{ix} , e_{iz} , e_{iz} are projections of unit vector e_i of spatial orientation n_i to the Cartesian coordinate axes x, y, z, separately. (i = 1, 2, 3, 4)

Proof. Assuming that i, j, k are the standard unit orthogonal vectors of the spatial coordinate system, based on the given conditions, one reads

$$e_i = e_{ix}i + e_{iy}j + e_{iz}k$$
, $i = 1, 2, 3, 4$

and

$$\begin{cases} e_{1x}\delta \boldsymbol{r} \cdot \boldsymbol{i} + e_{1y}\delta \boldsymbol{r} \cdot \boldsymbol{j} + e_{1z}\delta \boldsymbol{r} \cdot \boldsymbol{k} = \delta r_{1} \\ e_{2x}\delta \boldsymbol{r} \cdot \boldsymbol{i} + e_{2y}\delta \boldsymbol{r} \cdot \boldsymbol{j} + e_{2z}\delta \boldsymbol{r} \cdot \boldsymbol{k} = \delta r_{2} \\ e_{3x}\delta \boldsymbol{r} \cdot \boldsymbol{i} + e_{3y}\delta \boldsymbol{r} \cdot \boldsymbol{j} + e_{3z}\delta \boldsymbol{r} \cdot \boldsymbol{k} = \delta r_{3} \end{cases}$$

or

$$\begin{bmatrix} e_{1x} & e_{1y} & e_{1z} \\ e_{2x} & e_{2y} & e_{2z} \\ e_{3x} & e_{3y} & e_{3z} \end{bmatrix} \begin{pmatrix} \delta \boldsymbol{r} \cdot \boldsymbol{i} \\ \delta \boldsymbol{r} \cdot \boldsymbol{j} \\ \delta \boldsymbol{r} \cdot \boldsymbol{k} \end{bmatrix} = \begin{pmatrix} \delta r_1 \\ \delta r_2 \\ \delta r_3 \end{pmatrix}.$$

Because the three orientations e_1, e_2, e_3 are non-coplanar, the coefficient matrix is reversible. Hence, one can get

$$\begin{pmatrix} \delta \boldsymbol{r} \cdot \boldsymbol{i} \\ \delta \boldsymbol{r} \cdot \boldsymbol{j} \\ \delta \boldsymbol{r} \cdot \boldsymbol{k} \end{pmatrix} = \begin{bmatrix} e_{1x} & e_{1y} & e_{1z} \\ e_{2x} & e_{2y} & e_{2z} \\ e_{3x} & e_{3y} & e_{3z} \end{bmatrix}^{-1} \begin{pmatrix} \delta r_1 \\ \delta r_2 \\ \delta r_3 \end{pmatrix}.$$

Thus, the projection of the virtual displacement δr to orientation n_4 is

$$\delta \boldsymbol{r}_{4} = \left(\boldsymbol{e}_{4x}, \boldsymbol{e}_{4y}, \boldsymbol{e}_{4z}\right) \begin{pmatrix} \delta \boldsymbol{r} \cdot \boldsymbol{i} \\ \delta \boldsymbol{r} \cdot \boldsymbol{j} \\ \delta \boldsymbol{r} \cdot \boldsymbol{k} \end{pmatrix}$$
$$= \left(\boldsymbol{e}_{4x}, \boldsymbol{e}_{4y}, \boldsymbol{e}_{4z}\right) \begin{bmatrix} \boldsymbol{e}_{1x} & \boldsymbol{e}_{1y} & \boldsymbol{e}_{1z} \\ \boldsymbol{e}_{2x} & \boldsymbol{e}_{2y} & \boldsymbol{e}_{2z} \\ \boldsymbol{e}_{3x} & \boldsymbol{e}_{3y} & \boldsymbol{e}_{3z} \end{bmatrix}^{-1} \begin{pmatrix} \delta \boldsymbol{r}_{1} \\ \delta \boldsymbol{r}_{2} \\ \delta \boldsymbol{r}_{3} \end{pmatrix}.$$

Lemma 3. If the unit vectors e_i (i = 1, 2, 3, 4) are normal orthogonal unit vectors, the projection δr_4 of the virtual displacement to the orientation n_4 is

$$\delta r_{4} = \left(e_{4x}, e_{4y}, e_{4z}\right) \begin{bmatrix} e_{1x} & e_{2x} & e_{3x} \\ e_{1y} & e_{2y} & e_{3y} \\ e_{1z} & e_{2z} & e_{3z} \end{bmatrix} \begin{bmatrix} \delta r_{1} \\ \delta r_{2} \\ \delta r_{3} \end{bmatrix}.$$
(9)

Lemma 4. If the unit vectors e_i of orientations n_i (i = 1, 2, 3, 4) are non-planar vectors, and the projections of the virtual displacement of a dot to n_i are all zeros, then the virtual displacement must be zero vector, and projection of the virtual displacement to n_4 (every orientation) must be zero.

Lemma 5. If the unit vectors e_i (i = 1, 2, 3, 4) are standard unit orthogonal vectors of the Cartesian coordinate axes x, y, z, separately, the projection of the virtual displacement to e_i is δr_i , then the projection of the virtual displacement to the orientation n_4 is

$$\delta r_4 = e_{4x} \delta r_1 + e_{4y} \delta r_2 + e_{4z} \delta r_3.$$
 (10)

Example 3. Determine the force in member AC of the spatial truss, shown in Fig. 4 (a). Assume all members are pin connected. F_1 =5kN, F_2 =4kN, F_3 =2kN, F_4 =1.5kN, F_5 =1.5kN.

Solution. To calculate the force in rod AC, it is isolated from the system in Fig. 4 (b). Obviously, hinge points E, D, C and B are fixed. Therefore, the projection of point A to the orientations $B \rightarrow A$ and $D \rightarrow A$ are zeros. Assuming that the projection along the F_4 direction is δr , considering that the unit vectors of the direction $B \rightarrow A$, $A \rightarrow C$, $D \rightarrow A$ and F_4 are as follows

$$e_{BA} = \left(\frac{4}{5}, -\frac{3}{5}, 0\right), \quad e_{AC} = (-1, 0, 0),$$

 $e_{DA} = (0, 0, 1), \quad e_{F_4} = (0, -1, 0)$

Then, utilizing the second theorem for spatial case, the projection along the $A \rightarrow C$ direction is

$$\delta r_4 = (-1,0,0) \begin{bmatrix} \frac{4}{5} & -\frac{3}{5} & 0\\ 0 & 0 & 1\\ 0 & -1 & 0 \end{bmatrix}^{-1} \begin{pmatrix} 0\\ 0\\ \delta r \end{pmatrix}$$
$$= (-1,0,0) \begin{bmatrix} \frac{5}{4} & 0 & -\frac{3}{4}\\ 0 & 0 & -1\\ 0 & 1 & 0 \end{bmatrix} \begin{pmatrix} 0\\ 0\\ \delta r \end{pmatrix} = \frac{3}{4} \delta r$$

By employing the principle of virtual displacement, one reads

$$\boldsymbol{F}_{AC} \cdot \delta \boldsymbol{r}_A + \boldsymbol{F}_4 \cdot \delta \boldsymbol{r}_A = 0. \tag{11}$$

Based on the above analysis and (11), one can get

$$F_{AC} \cdot \frac{3}{4} \delta r + F_4 \cdot \delta r = 0$$

Because $\delta \theta \neq 0$, then one can get

$$F_{AC} = -\frac{4}{3}F_4 = -\frac{4}{3} \times 1.5$$
kN = -2.0kN.



Fig. 4 Example 3

4 Conclusions

This paper mainly faces to the difficulty of computation of virtual work, investigates on the projection of a dot's virtual displacement to a given orientation, propose two theorems and corresponding lemmas, and discuss their application for analyzing the forces of structural members. Computation procession reveals that the formula for planar situation is easy to use, and the formula for spatial situation is normalized and easy to be remembered. Combining with current methods, these results will be very helpful to analyze inner forces of complicit structures. It's necessary to say that the two theorems can be used in dynamics when dynamical problems convert into static problems in form by utilizing D' Alembert's principle.

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