

QUASI-RECURRENT WEYL SPACES

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Abstract

In this work we define Quasi-Recurrent Weyl spaces and examine the hypersurfaces of them.

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1. Introduction

An n -dimensional manifold W_n is said to be a Weyl space if it has a conformal metric tensor g_{ij} and a symmetric connection ∇_k satisfying the compatibility condition given by the equation

$$\nabla_k g_{ij} - 2T_k g_{ij} = 0, \quad (1)$$

where T_k denotes a covariant vector field and $\nabla_k g_{ij}$ denotes the usual covariant derivative.

Under renormalization of the fundamental tensor of the form

$$\check{g}_{ij} = \lambda^2 g_{ij} \quad (2)$$

the complementary vector T_k is transformed by the law

$$\check{T}_i = T_i + \partial_i \ln \lambda, \quad (3)$$

where λ is a function.

Writing (1) out in full, we have

$$\partial_k g_{ij} - g_{hj} \Gamma_{ik}^h - g_{ih} \Gamma_{jk}^h - 2T_k g_{ij} = 0,$$

where Γ_{kl}^i are the connection coefficients of the form

$$\Gamma_{kl}^i = \left\{ \begin{matrix} i \\ kl \end{matrix} \right\} - g^{im} (g_{mk} T_l + g_{ml} T_k - g_{kl} T_m). \quad (4)$$

A quantity A is called a satellite with weight $\{p\}$ of the tensor g_{ij} , if it admits a transformation of the form

$$\check{A} = \lambda^p A$$

under the renormalization (2) of the metric tensor g_{ij} .

The prolonged covariant derivative of a satellite A of the tensor g_{ij} with weight $\{p\}$ is defined by [1]

$$\check{\nabla}_k A = \nabla_k A - p T_k A. \quad (5)$$

REMARK 1 The prolonged covariant derivative preserves the weight.

The curvature tensor R_{ijkl}^i of the Weyl connection is defined by

$$R_{jkl}^i = \frac{\partial}{\partial x^k} \Gamma_{jl}^i - \frac{\partial}{\partial x^l} \Gamma_{jk}^i + \Gamma_{hk}^i \Gamma_{jl}^h - \Gamma_{hl}^i \Gamma_{jk}^h \quad (6)$$

and the Ricci tensor R_{ij} of the Weyl connection is

$$R_{ij} = R_{ijm}^m. \quad (7)$$

Since the Weyl connection is not metric, the Ricci tensor R_{ij} is not necessarily symmetric. In fact, $R_{[ij]} = n \nabla_{[i} T_{j]}$. We remark that if T_j is a gradient, then the space is Riemannian.

It is easy to see that the covariant curvature tensor R_{lij}^k is a satellite of g_{ij} with weight $\{2\}$.

The Bianchi identity for the Weyl space is, by [2]

$$\check{\nabla}_l R_{ijk}^h + \check{\nabla}_k R_{lij}^h + \check{\nabla}_j R_{ikl}^h = 0. \quad (8)$$

2. Quasi-Recurrent Weyl Spaces

A non-flat Weyl space $W_n(g_{ij}, T_k)$ will be called quasi-recurrent ($(QRW)_n$ in short) if the curvature tensor satisfies the following condition for some non-zero covariant vector field ϕ_k ($\neq T_k$)

$$\check{\nabla}_s R_{ijkl} = 2\phi_s R_{ijkl} + \phi_i R_{sjkl} + \phi_j R_{iskl} + \phi_k R_{ijsl} + \phi_l R_{ijks}. \quad (9)$$

A $(QRW)_n$ manifold can be Weyl recurrent, i.e. it can, beside (9), satisfy

$$\check{\nabla}_s R_{ijhk} = \phi'_s R_{ijhk}. \quad (10)$$

We examine the spaces satisfying (9) but not satisfying (10). Recurrent Weyl spaces have been examined in [2]. We note that ϕ_k is a satellite of g_{ij}

with weight $\{0\}$. By multiplying (9) by g^{il} and summing up with respect to i and l we get

$$\dot{\nabla}_s R_{jk} = 2\phi_s R_{jk} + \phi_j R_{sk} + \phi_k R_{js} + o_i(R_{sjk}^i + R_{jk_s}^i). \quad (11)$$

Similarly transvecting (11) by g^{jk} we obtain

$$\dot{\nabla}_s R = 2\phi_s R + \phi_k (R_s^k + R_{.s}^k + R_{si..}^{ik} + R_{.is}^{ki}), \quad (12)$$

where $R_s^k = R_{sj}g^{kj}$ and $R_{.s}^k = R_{js}g^{jk}$.

Hence the scalar curvature of the $(QRW)_n$ satisfies (12).

By changing the indices j and k in (11) and subtracting the resulting equation from the one obtained therefrom we obtain

$$\dot{\nabla}_s F_{jk} = \frac{2(n+1)}{n} \phi_s F_{jk} + \phi_k F_{js} + \phi_j F_{sk}, \quad (13)$$

where $F_{jk} = R_{[jk]}$.

In fact this is the relation between the complementary vector T_k and the recurrent vector ϕ_k .

3. Hypersurfaces of Quasi-Recurrent Weyl Spaces

Let $W_n(g_{ij}, T_k)$ be a hypersurface, with coordinates u^i ($i = 1, 2, \dots, n$) of a Weyl space $W_{n+1}(g_{ab}, T_c)$ with coordinates x^a ($a = 1, 2, \dots, n+1$). The metrics of W_n and W_{n+1} are connected by the relations

$$g_{ij} = g_{ab} x_i^a x_j^b \quad (i, j = 1, 2, \dots, n; \quad a, b = 1, 2, \dots, n+1). \quad (14)$$

where x_i^a denotes the covariant derivative of x^a with respect to u^i .

It is easy to see that the prolonged covariant derivative of a satellite A , relative to W_n and W_{n+1} , are related by

$$\dot{\nabla}_k A = x_k^c \dot{\nabla}_c A \quad (k = 1, 2, \dots, n; \quad c = 1, 2, \dots, n+1). \quad (15)$$

Let n^a be the contravariant components of the vector field of W_{n+1} normal to W_n which is normalized by the condition

$$g_{ab} n^a n^b = 1. \quad (16)$$

The moving frame $\{x_a^i, n_a\}$ in W_n , reciprocal to the moving frame $\{x_i^a, n^a\}$ is defined by the relations

$$n_a x_i^a = 0, \quad n^a x_a^i = 0, \quad x_i^a x_a^j = \delta_i^j. \quad (17)$$

Remembering that the weight of x_i^a is $\{0\}$, the prolonged covariant derivative of x_i^a with respect to u^k is found as

$$\dot{\nabla}_k x_i^a = \nabla_k x_i^a = \omega_{ik} n^a, \quad (18)$$

where ω_{ik} is the second fundamental form. It can be shown that ω_{ik} is a satellite of g_{ij} with weight $\{1\}$.

The following two relations, which are respectively the generalization of Gauss and Mainardi-Codazzi equations, are obtained in [2]

$$R_{pijk} = \Omega_{pijk} + \bar{R}_{dbce} x_p^d x_i^b x_j^c x_k^e, \quad (19)$$

$$\dot{\nabla}_k \omega_{ij} - \dot{\nabla}_j \omega_{ik} + \bar{R}_{dbce} x_i^b x_j^c x_k^e n^d = 0, \quad (20)$$

where \bar{R}_{dbce} is the covariant curvature tensor of W_{n+1} and Ω_{pijk} is the Sylvesterian of ω_{ij} defined by $\Omega_{pijk} = \omega_{pj} \omega_{ik} - \omega_{pk} \omega_{ij}$.

THEOREM 1 A hypersurface W_n of $(QRW)_{n+1}$ satisfies the following.

$$\begin{aligned} \dot{\nabla}_s R_{ijkl} - \dot{\nabla}_s \Omega_{ijkl} &= 2\phi_s(R_{ijkl} - \Omega_{ijkl}) + \phi_i(R_{sjkl} - \Omega_{sjkl}) + \\ &+ \phi_j(R_{iskl} - \Omega_{iskl}) + \phi_k(R_{ijsl} - \Omega_{ijsl}) + \\ &+ \phi_l(R_{ijks} - \Omega_{ijks}) + \bar{R}_{abcd} \dot{\nabla}_s(x_i^a x_j^b x_k^c x_l^d). \end{aligned} \quad (21)$$

Proof 1 By taking the prolonged covariant derivative of (19) we have

$$\dot{\nabla}_s R_{ijkl} = \dot{\nabla}_s \Omega_{ijkl} + \dot{\nabla}_s(\bar{R}_{abcd}) x_i^a x_j^b x_k^c x_l^d + \bar{R}_{abcd} \dot{\nabla}_s(x_i^a x_j^b x_k^c x_l^d).$$

Moreover,

$$\dot{\nabla}_s R_{ijkl} = \dot{\nabla}_s \Omega_{ijkl} + \dot{\nabla}_e(\bar{R}_{abcd}) x_i^a x_j^b x_k^c x_s^e + \bar{R}_{abcd} \dot{\nabla}_s(x_i^a x_j^b x_k^c x_l^d).$$

By using (9), Mainardi-Codazzi equations and (15), we obtain the result.

THEOREM 2 A simply connected hypersurface of $(QRW)_{n+1}$ is Riemannian.

Proof 2 If we change the indices k, l, s cyclically in (21), we obtain two more equations. Namely

$$\begin{aligned} \dot{\nabla}_k R_{ijls} - \dot{\nabla}_k \Omega_{ijls} &= 2\phi_k(R_{ijls} - \Omega_{ijls}) + \phi_i(R_{kjls} - \Omega_{kjls}) + \\ &+ \phi_j(R_{ikls} - \Omega_{ikls}) + \phi_l(R_{ijks} - \Omega_{ijks}) + \\ &+ \phi_s(R_{ijlk} - \Omega_{ijlk}) + \bar{R}_{abcd} \dot{\nabla}_k(x_i^a x_j^b x_l^c x_s^d), \end{aligned} \quad (22)$$

$$\begin{aligned} \dot{\nabla}_l R_{ijsk} - \dot{\nabla}_l \Omega_{ijsk} &= 2\phi_l(R_{ijsk} - \Omega_{ijsk}) + \phi_i(R_{ljsk} - \Omega_{ljsk}) + \\ &+ \phi_j(R_{ilsk} - \Omega_{ilsk}) + \phi_s(R_{ijlk} - \Omega_{ijlk}) + \\ &+ \phi_k(R_{ijsl} - \Omega_{ijsl}) + \bar{R}_{abcd} \dot{\nabla}_l(x_i^a x_j^b x_s^c x_k^d). \end{aligned} \quad (23)$$

Adding these three equations and using Mainardi-Codazzi equations and Bianchi identities, we find

$$\phi_i(R_{lj sk} + R_{kj ls} + R_{sj kl}) = 0. \quad (24)$$

Since $\phi_i \neq 0$ we have

$$R_{lj sk} + R_{kj ls} + R_{sj kl} = 0. \quad (25)$$

From the first Bianchi identity we get

$$R_{(lj)sk} + R_{(kj)ls} + R_{(sj)kl} = 0. \quad (26)$$

Hence, the result follows from the fact that $R_{(ij)sk} = 2g_{ij}T_{[s,k]}$.

A hypersurface W_n of W_{n+1} is called totally geodesic if $\omega_{ij} = 0$. Therefore we have the following theorem.

THEOREM 3 If the simply connected hypersurface of $(QRW)_{n+1}$ is totally geodesic, then the hypersurface is Riemannian quasi-recurrent.

Proof 3 Follows from (21) and the Theorem 2.

References

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